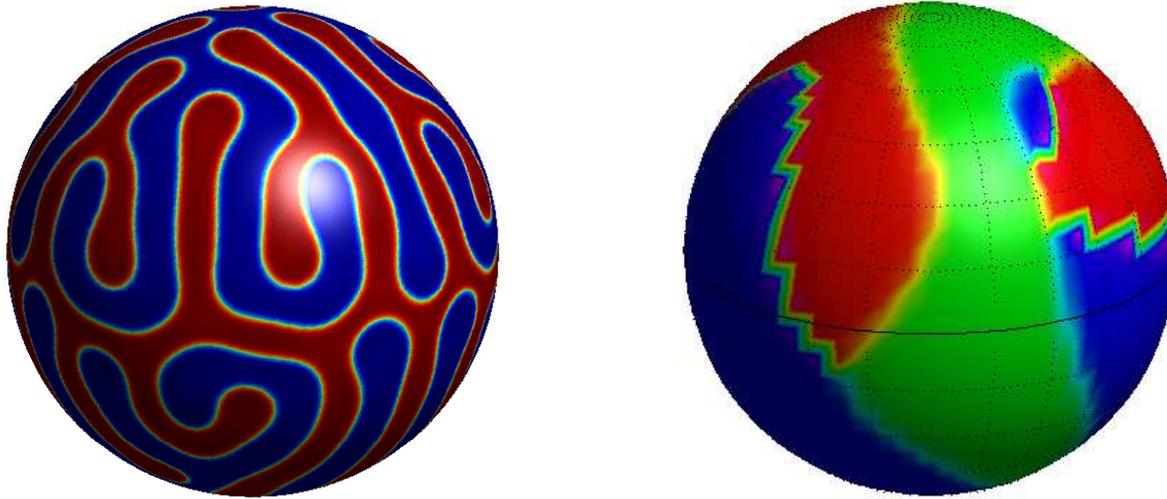
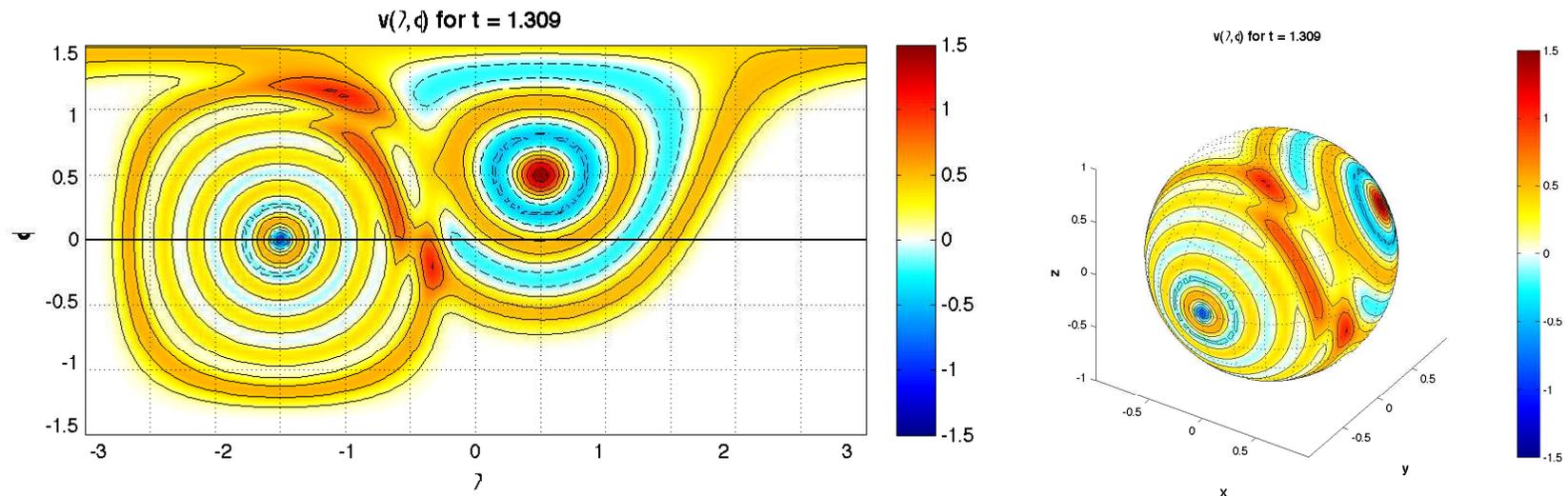


# Fourier Spectral Computing for PDEs on the Sphere

- ▷ an FFT-based method with implicit-explicit timestepping
- ▷ a simple & efficient approach



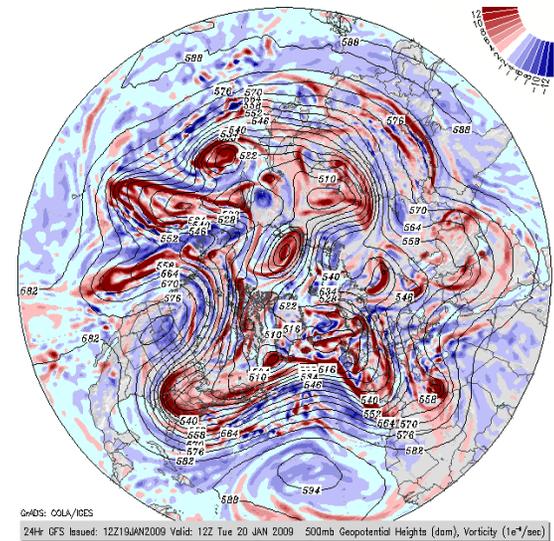
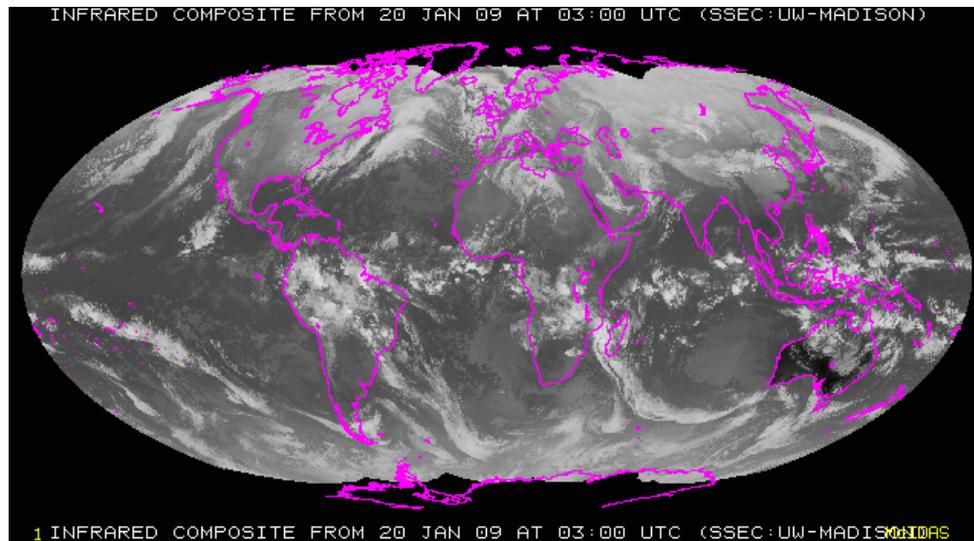
- ▷ Dave Muraki, Andrea Blazenko & Kevin Mitchell  
Mathematics, Simon Fraser University



## Remote Participation

- ▷ simple Matlab demos for familiar PDEs can be run in real-time on laptop
  - ▷ diffusion & wave equations
  - ▷ reaction-diffusion & phase field patterns
  - ▷ nonlinear Schrödinger equation
- ▷ operation count: DFTs on sphere scale as 2D FFTs
  - ▷ all latitude calculations done in  $O(N)$  operations: scaling reduced by factor of  $N \sim 64$ -256
- ▷ matlab demos indicated by footnotes: new demo indicated in red

# PDEs on the Sphere: Applications & Computation

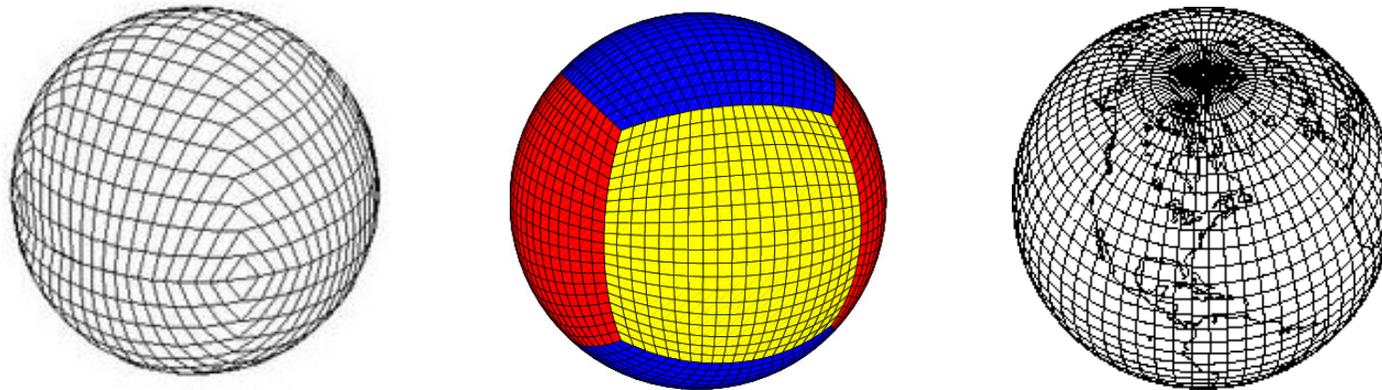


## Applications

- ▷ computing PDEs on the sphere driven by geophysics
  - ▷ meteorology: atmospheric fluid dynamics
  - ▷ climatology: aqua-planet oceanography
  - ▷ seismology: Rayleigh surface waves
- ▷ Fourier analysis on the sphere
  - ▷ tomography, crystallography, computer graphics
- ▷ computing on a manifold: constant, positive curvature

# PDEs on the Sphere: Applications & Computation

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## Computation

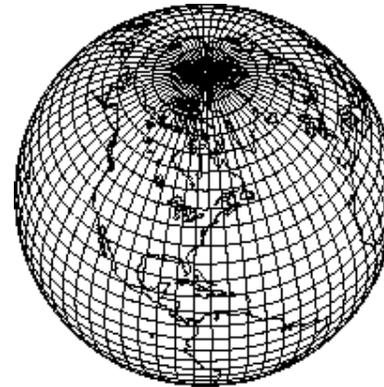
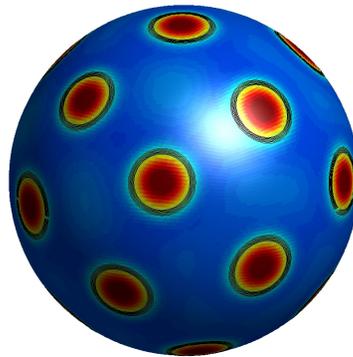
- ▷ numerical schemes
  - ▷ finite-difference, finite-volume, finite-element, spectral element . . .
  - ▷ spherical harmonics
- ▷ gridding
  - ▷ logically-rectangular, cubed sphere, longitude-latitude (long-lat), *yin-yang* overset grid . . .
- ▷ parallelization, mesh refinement & adaptivity

## Fourier Spectral Method

- ▷ spectrally-fast: uses FFT for Fourier-based spectral transform
- ▷ simple implementation: uses long-lat grid & resembles FFT computing on 2D periodic rectangle

# PDEs on the Sphere: Applications & Computation

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## Computation

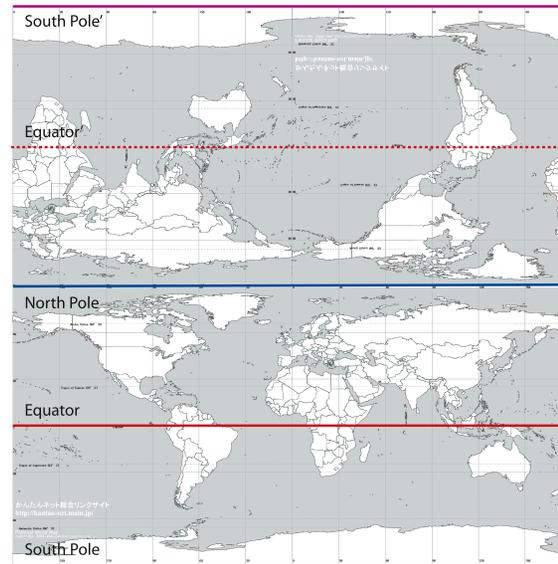
- ▷ numerical schemes [Calhoun, Helzel & LeVeque, 2008]
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# Periodicity of the Sphere

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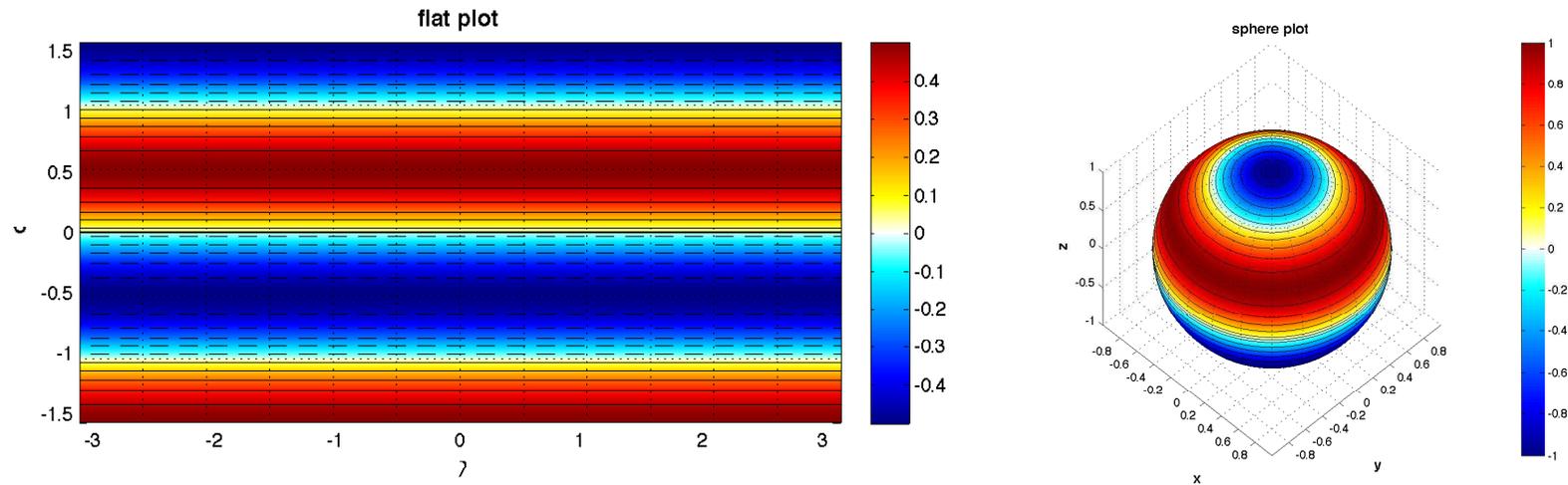
When is a Sphere is not a Sphere? . . . When it is a Torus!

- ▷ double mapping of a sphere to a periodic rectangle
  - ▷ north (NP) & south pole (SP) become lines of constant value
  - ▷ longitude  $\rightarrow \lambda$ -axis ( $-\pi \leq \lambda \leq +\pi$ )
  - ▷ co-latitude  $\rightarrow \phi$ -axis ( $0 \leq \phi \leq 2\pi$ )
- ▷ spherical symmetry: smooth extension from long-lat sphere to torus

$$f(\lambda, \phi) = \begin{cases} f(\lambda, \phi) & \text{for } 0 \leq \phi \leq \pi \\ f(\pi - \lambda, 2\pi - \phi) & \text{for } \pi \leq \phi \leq 2\pi \end{cases} .$$

- ▷ PDE should preserve spherical symmetry

# Fourier Modes on the Sphere



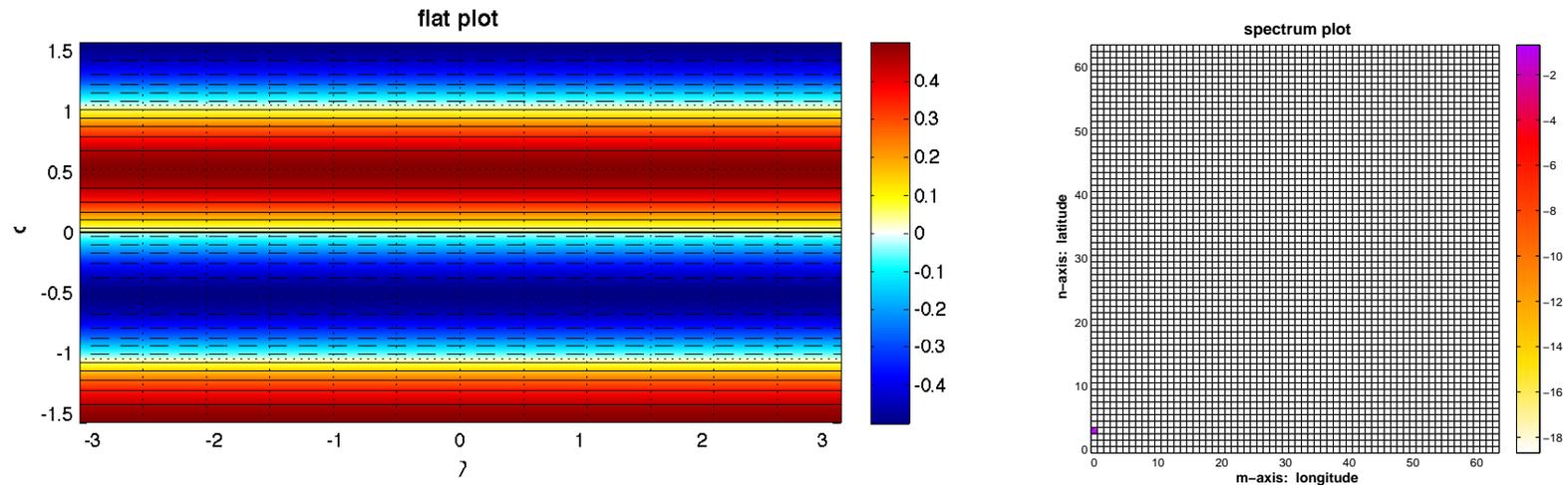
## Fourier Modes with Spherical Symmetry ( $m = 0$ )

- ▷ natural Fourier modes in longitude:  $Q_{nm}(\lambda, \phi) = q_{nm}(\phi) e^{im\lambda}$
- ▷ trigonometric modes for latitude:

$$q_{nm}(\phi) = \begin{cases} \cos n\phi & \text{for } m = 0 \\ \sin \phi \sin n\phi & \text{for } m \text{ even} \\ \sin n\phi & \text{for } m \text{ odd} \end{cases}$$

- ▷ smoothness at the poles
  - ▷  $m$  even modes have even symmetry across poles & are zero at poles for  $m \neq 0$
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# Fourier Modes on the Sphere



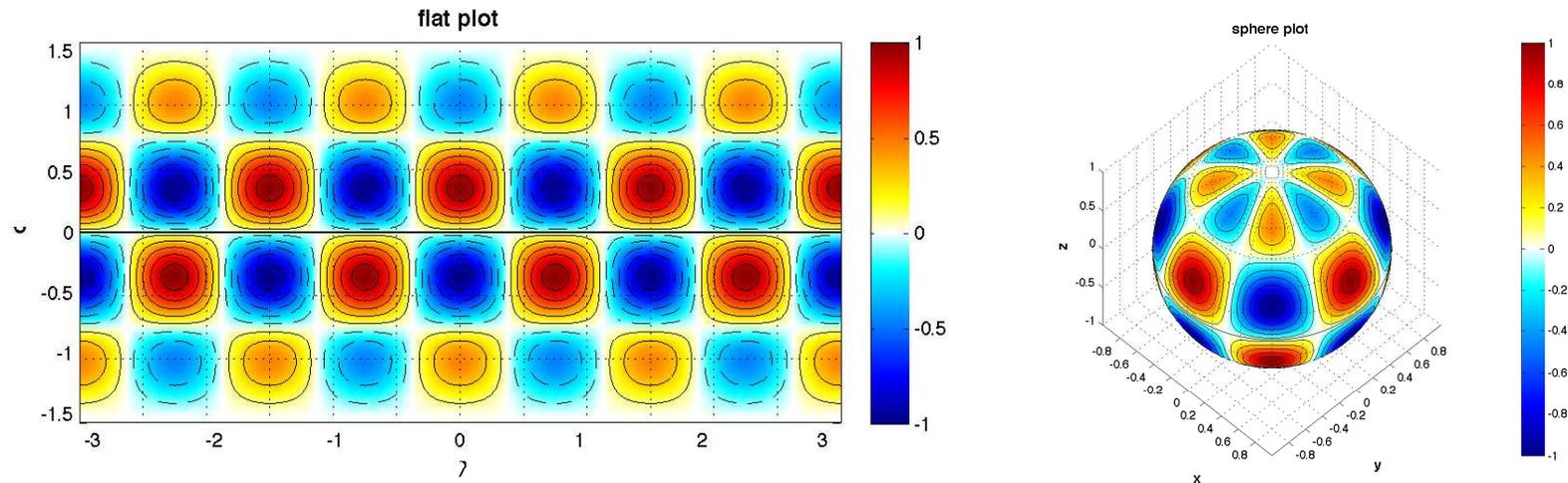
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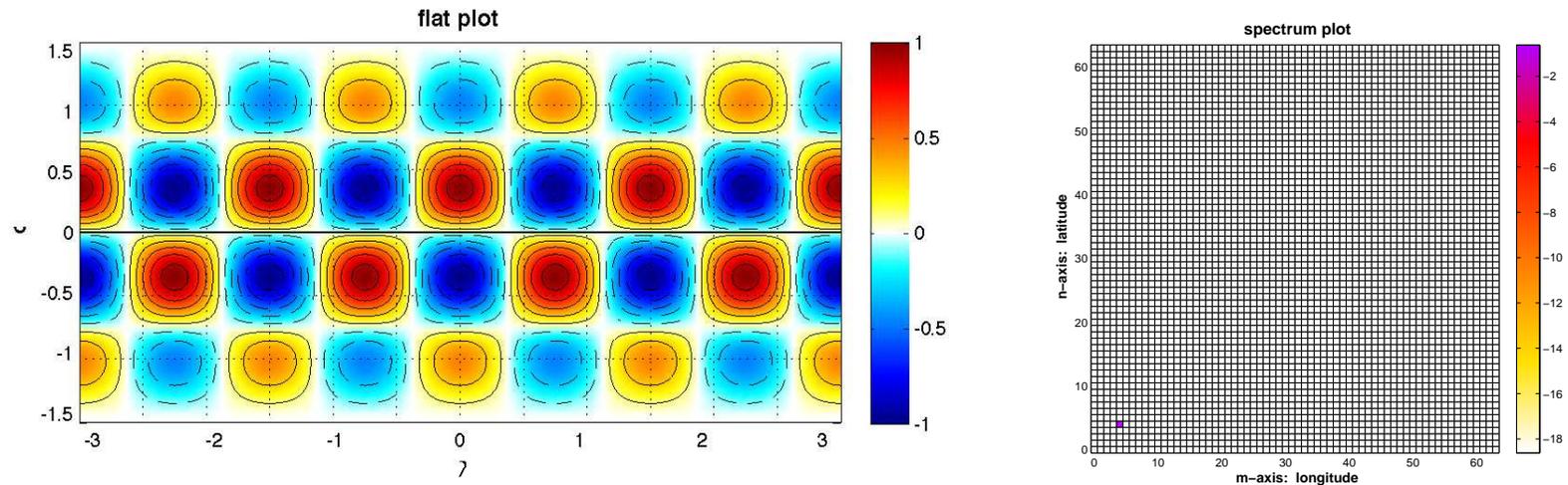
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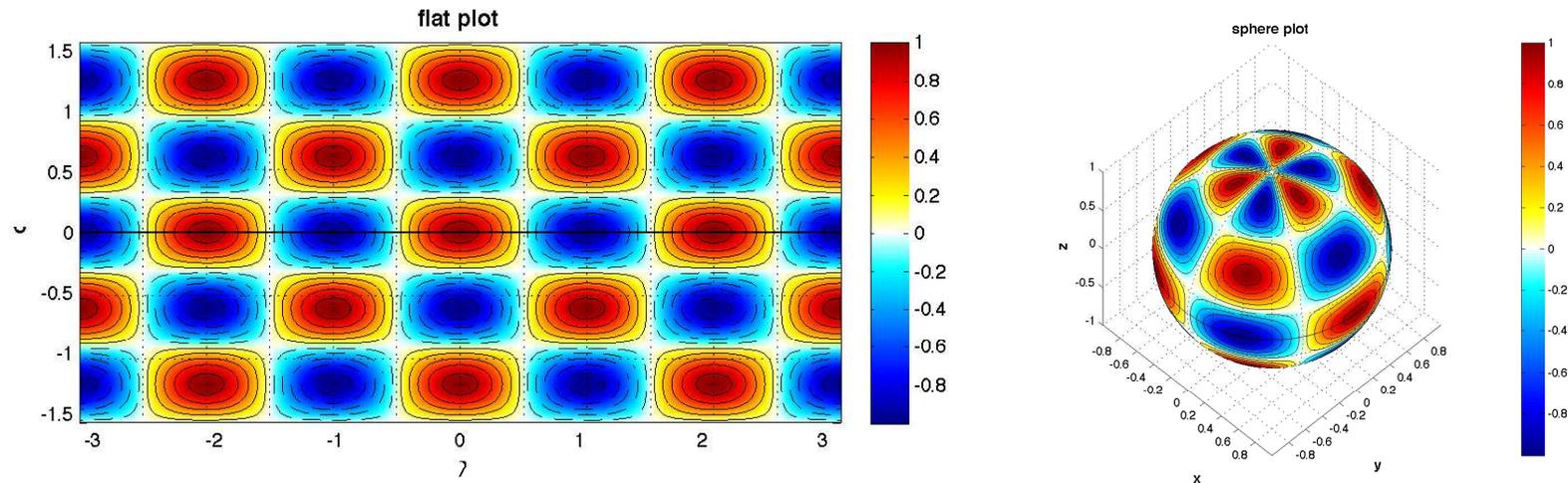
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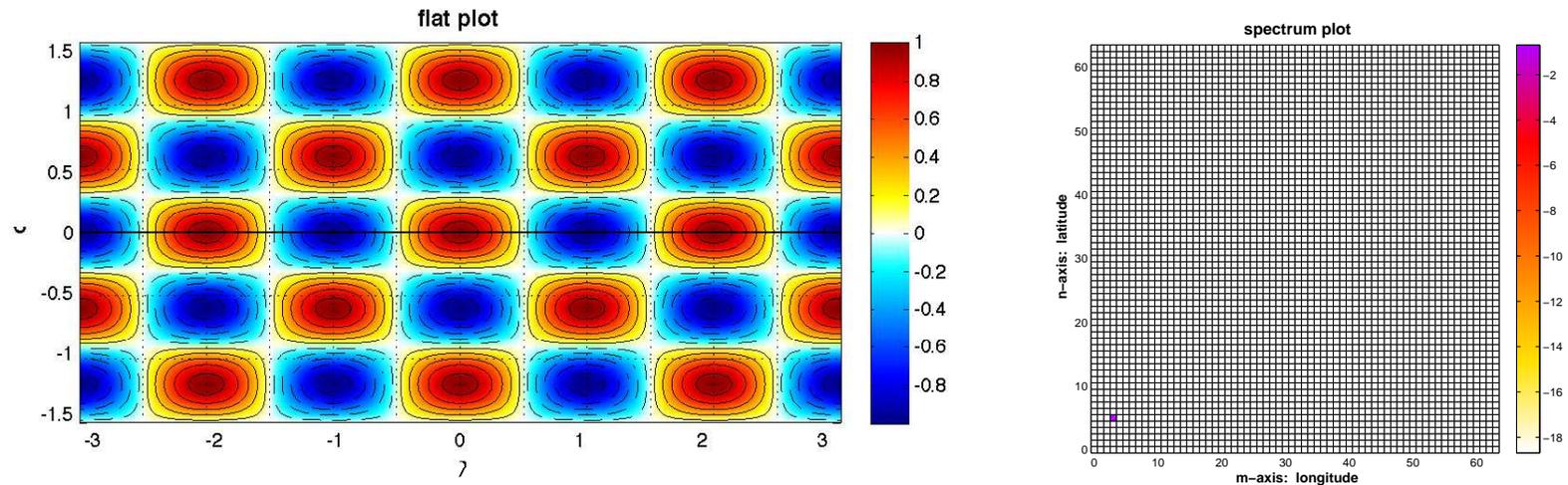
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# Fourier Modes on the Sphere



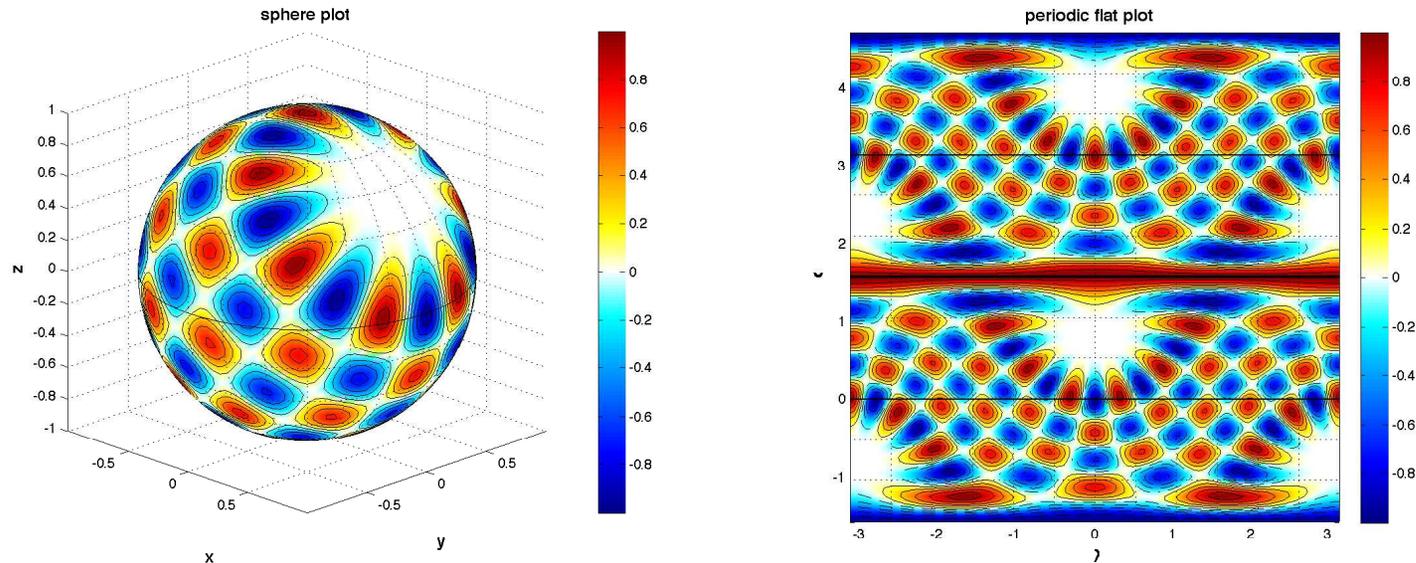
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# Double Fourier Series on the Sphere



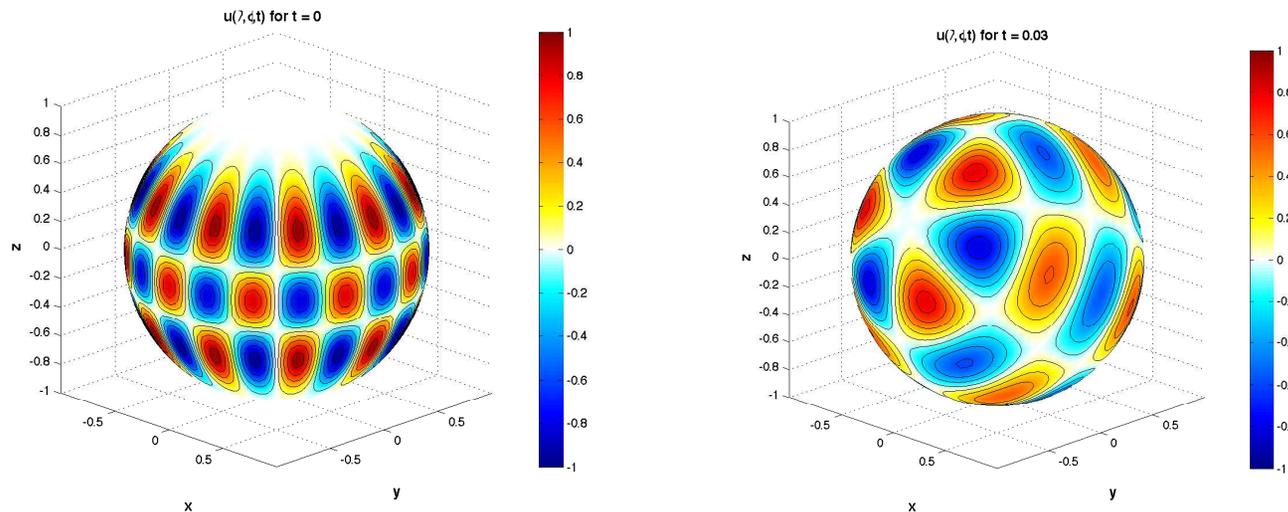
## Fourier Geometry for Spherically-Symmetric Functions

- ▷ double sum over  $Q_{nm}$  modes

$$f(\lambda, \phi, t) = \sum_{m=-\infty}^{+\infty} \sum_{n=0}^{\infty} \tilde{f}_{nm}(t) Q_{nm}(\lambda, \phi) = \sum_{m=-\infty}^{+\infty} e^{im\lambda} \sum_{n=0}^{\infty} \tilde{f}_{nm}(t) q_{nm}(\phi)$$

- ▷ early works: Merilees, (1973), Orszag (1973), Boer/Steinberg (1974), Boyd (1978)
- ▷ elliptic solves: Yee (1981), Moorthi/Higgins (1992)
- ▷ fluid flow: Fornberg (1997), Spatz, et.al. (1998), Cheong et.al. (2000-2006), Layton/Spatz (2003)
- ▷ PDE must respect spherical symmetry

# Diffusion on the Sphere



## Initial Decay and Steady-State Forcing, $N = 64$

- ▷ forced diffusion on sphere,  $u(\lambda, \phi, t)$

$$u_t = \nabla^2 u + f(\lambda, \phi)$$

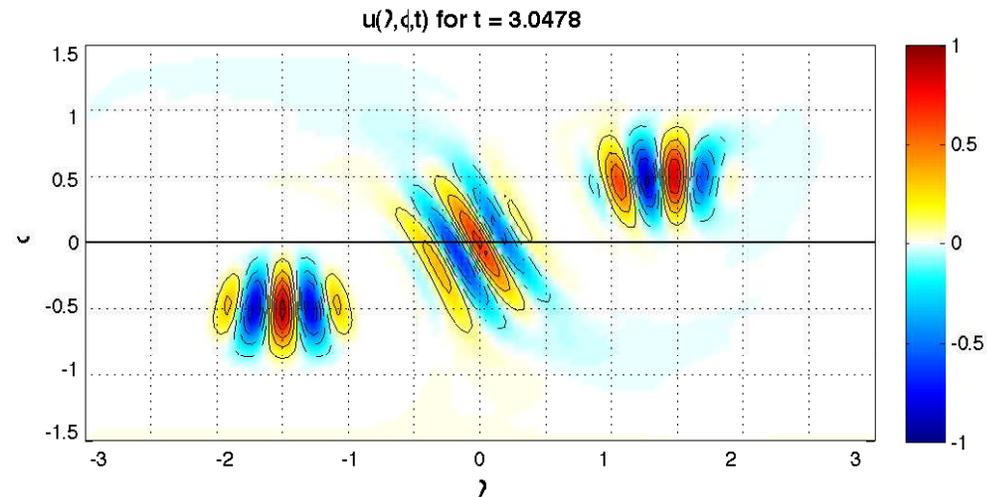
- ▷ initial condition:  $u_0(\lambda, \phi, 0) = c_1 \{Y_{12}^{10}(\lambda, \phi) + \text{conj}\}$
- ▷ steady forcing,  $f(\lambda, \phi)$  is rotated spherical harmonic:  $c_2 \{Y_7^3(\lambda, \phi) + \text{conj}\}$
- ▷ exponential-in- $t$  solutions
$$u(\lambda, \phi, t) = e^{-156t} u_0(\lambda, \phi) + (1 - e^{-72t}) f(\lambda, \phi)$$

- ▷ FFT-based routines for spherical data: *fftS.m* & *ifftS.m*

- ▷ 3<sup>rd</sup>-order accurate timestepping

# Waves on the Sphere

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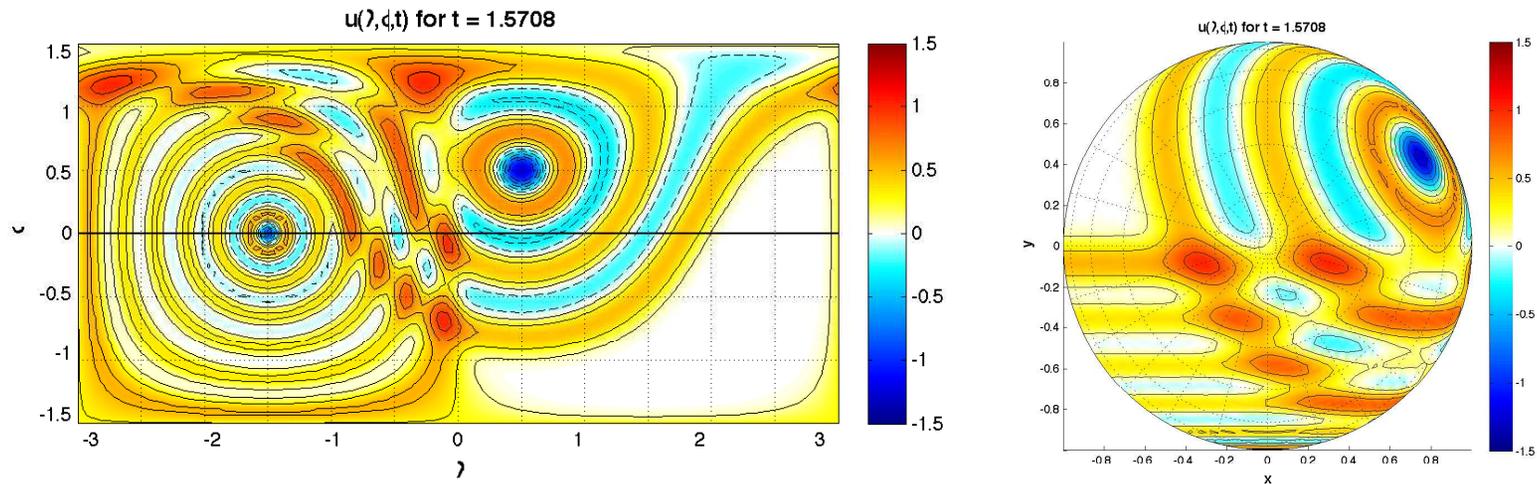
## Uni-Directional Wavepacket, $N = 64$

- ▷ propagation on sphere,  $u(\lambda, \phi, t)$  &  $v(\lambda, \phi, t)$

$$\begin{aligned} u_t &= v \\ v_t &= \nabla^2 u \end{aligned}$$

- ▷ initial wavepacket: masked spherical harmonic,  $u_0(\lambda, \phi) = c_1 \{Y_{16}^{16}(\lambda, \phi) + \text{conj}\}$
- ▷ propagation follows a great circle
- ▷ wavespeeds  $\rightarrow 1^+$ , in short wave limit
- ▷ 2<sup>nd</sup>-order accurate timestepping scheme for stability

# Waves on the Sphere



## Forced Waves, $N = 64$

- ▷ radiation on the sphere,  $u(\lambda, \phi, t)$  &  $v(\lambda, \phi, t)$

$$\begin{aligned}u_t &= v \\v_t &= \nabla^2 u + f_1(\lambda, \phi) \sin 24t + f_2(\lambda, \phi) \sin 12t\end{aligned}$$

- ▷ zero initial condition:  $u_0(\lambda, \phi, 0) = 0$
- ▷ oscillatory forcings:  $f_j(\lambda, \phi)$  are spatially localized
- ▷ 2<sup>nd</sup>-order accurate timestepping
- ▷ no problems at the pole: geometrical distortions, or time-step restriction from over-resolution

# Spectral Computing on the Sphere

---

## Spherical Harmonics (SH)

- ▷ characteristics of  $Y_l^m(\lambda, \phi)$ 
  - ▷  $l$ -index indicates spatial resolution on the sphere
  - ▷ SH modes preserve spatial resolution under coordinate rotations
  - ▷ SH spectrum is Fourier in  $\lambda$ -direction;  $m$ -index gives direction ( $-l \leq m \leq +l$ )
  - ▷ eigenfunctions of the surface Laplacian
- ▷ fast SH transform (S2kit), based on Driscoll & Healy (1989)
  - ▷ high-complexity algorithm: fast interpolations, fast multipole, WKB approximation . . .
  - ▷ Orszag (1986), Jakob-Chien/Alpert (1997), Suda/Takami (2001)

## PDE Computing with Spherical Harmonics

- ▷ spectral method of choice for geophysical fluid codes
- ▷ elliptic solves & timestepping schemes
- ▷ Swartztrauber's 1979 assessment:

*“the theoretical gap which exists between the states of the art for discrete spectral approximations on a sphere and on a rectangle.”*

# Spatial Derivatives on the Sphere

---

Spectral Differentiation of  $Q_{nm}(\lambda, \phi) = q_{nm}(\phi) e^{im\lambda}$

- ▷ longitude differentiation

$$\frac{\partial}{\partial \lambda} Q_{nm} = im Q_{nm}$$

- ▷ latitude differentiation, typically non-constant coefficient

$$\sin \phi \frac{\partial}{\partial \phi} Q_{nm} = a_1 Q_{(n-2)m} + a_2 Q_{nm} + a_3 Q_{(n+2)m}$$

- ▷ **tri-diagonal** differentiation matrices:  $\sin^2 \phi \nabla^2$  ,  $\sin^2 \phi$
- ▷ forward & inverse operations remain  $O(N^2)$

## Double Fourier Series

- ▷ derivative operations on the spectral representation . . .

$$u(\lambda, \phi, t) = \sum_m e^{im\lambda} \sum_n \tilde{u}_{nm}(t) q_{nm}(\phi)$$

. . . act on Fourier coefficients  $\tilde{u}_{nm}$  as vectors  $(\tilde{u}_n)_m$

# Timestepping for Fourier Methods

---

## Diffusion on a Rectangle

- ▷ Fourier representation

$$u(x, y, t) = \sum_m \sum_n \tilde{u}_{nm}(t) e^{i(mx+ny)}$$

- ▷ forced diffusion on sphere

$$u_t - \nabla^2 u = f(\lambda, \phi, t)$$

- ▷ spectral equation

$$\frac{d}{dt} \tilde{u}_{nm} + (m^2 + n^2) \tilde{u}_{nm} = \tilde{f}_{nm}(t)$$

## Fourier Timestepping Strategies

- ▷ **integrating factor** treats Laplacian exactly  $\rightarrow$  ODE solve

$$\frac{d}{dt} \left\{ e^{(m^2+n^2)t} \tilde{u}_{nm} \right\} = e^{(m^2+n^2)t} \tilde{f}_{nm}(t)$$

- ▷ exponential time-differencing  $\rightarrow$  numerical **quadrature**

$$\tilde{u}_{nm}(t) = \tilde{u}_{nm}(0) e^{-(m^2+n^2)t} + \int_0^t e^{(m^2+n^2)(s-t)} \tilde{f}_{nm}(s) ds$$

# IMEX Timestepping for Fourier on Sphere

---

## Diffusion on the Sphere

- ▷ Fourier representation &  $(\tilde{u}_n)_m$  as column vector data

$$u(x, y, t) = \sum_m \sum_n (\tilde{u}_n)_m q_{nm}(\phi) e^{im\lambda}$$

- ▷ forced diffusion on sphere

$$\sin^2 \phi u_t - \sin^2 \phi \nabla^2 u = \sin^2 \phi f(\lambda, \phi, t)$$

- ▷ spectral equation

$$[\sin^2 \phi] \frac{d}{dt} (\tilde{u}_n)_m - [\sin^2 \phi \nabla^2] (\tilde{u}_n)_m = [\sin^2 \phi] (\tilde{f}_n)_m$$

## IMEX Scheme [Ascher, Ruuth, Wetton (1995)]

- ▷ 2<sup>rd</sup>-order BDF in time

$$[\sin^2 \phi] \left\{ \frac{3(\tilde{u}_n)_m^{j+1} - 4(\tilde{u}_n)_m^j + (\tilde{u}_n)_m^{j-1}}{2\Delta t} \right\} - [\sin^2 \phi \nabla^2] (\tilde{u}_n)_m^{j+1} = [\sin^2 \phi] \left\{ 2(\tilde{f}_n)_m^j - (\tilde{u}_n)_m^{j-1} \right\}$$

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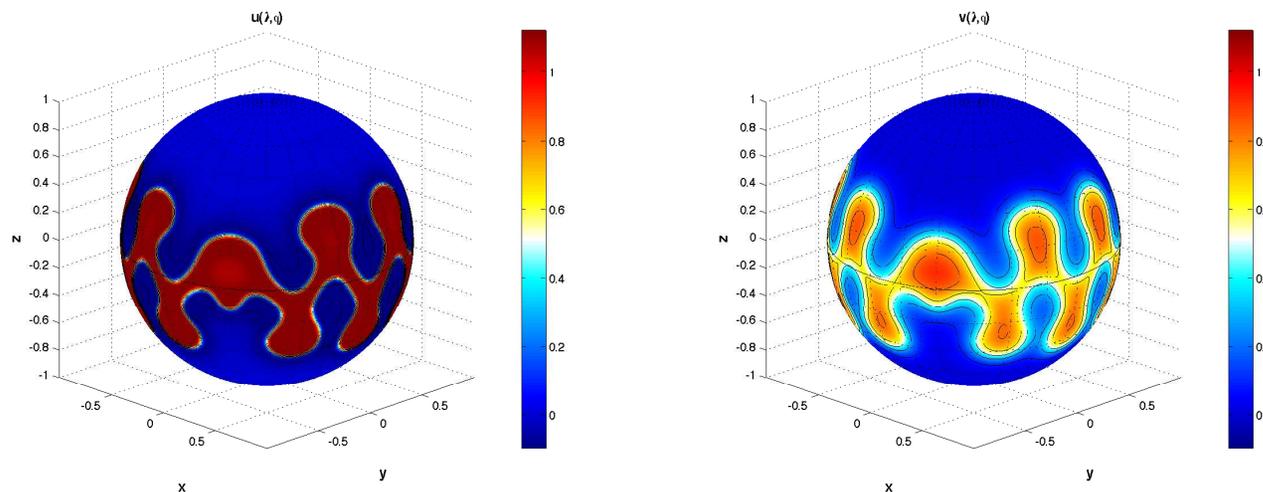
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## IMEX Scheme [Ascher, Ruuth, Wetton (1995)]

- ▷ 2<sup>rd</sup>-order BDF in time

$$\left\{ \frac{3}{2\Delta t} [\sin^2 \phi] - [\sin^2 \phi \nabla^2] \right\} (\tilde{u}_n)_m^{j+1} =$$
$$[\sin^2 \phi] \left\{ \frac{4(\tilde{u}_n)_m^j - (\tilde{u}_n)_m^{j-1}}{2\Delta t} + 2(\tilde{f}_n)_m^j - (\tilde{u}_n)_m^{j-1} \right\}$$

# Application: Fitz-Hugh Nagumo Equations



## Pattern Formation & Front Dynamics, $N = 256$

- ▷ bi-stable  $(0, 1)$  activator,  $u(\lambda, \phi, t)$  & long-range inhibitor,  $v(\lambda, \phi, t)$

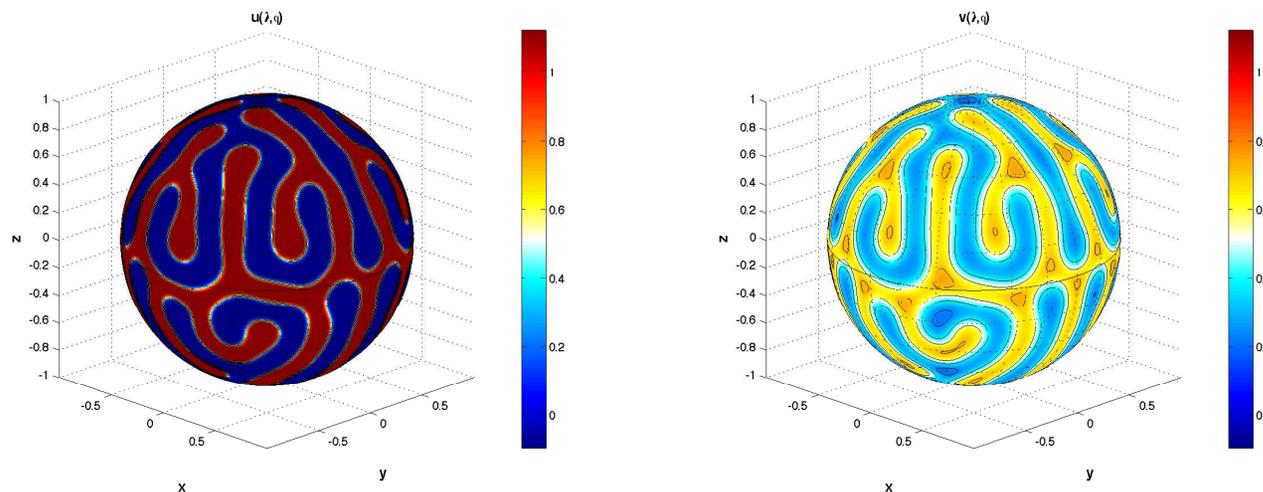
$$u_t = \frac{\epsilon^2}{r^2} \nabla^2 u + u(u - a)(u - 1) + \rho(v - u)$$

$$0 = \frac{1}{r^2} \nabla^2 v + (v - u)$$

- ▷ labyrinth forming region in parameter space [Goldstein, DJM, Petrich, 1992]

- ▷  $0 < \epsilon \ll 1 \ll r$  → sharp fronts in  $u$  & large spherical domain
- ▷  $0 < a - 1/2 \ll 1$  →  $u = 0, 1$  bistability with weak bias to red
- ▷  $0 < \rho \ll 1$  → weak inhibition producing blue

# Application: Fitz-Hugh Nagumo Equations



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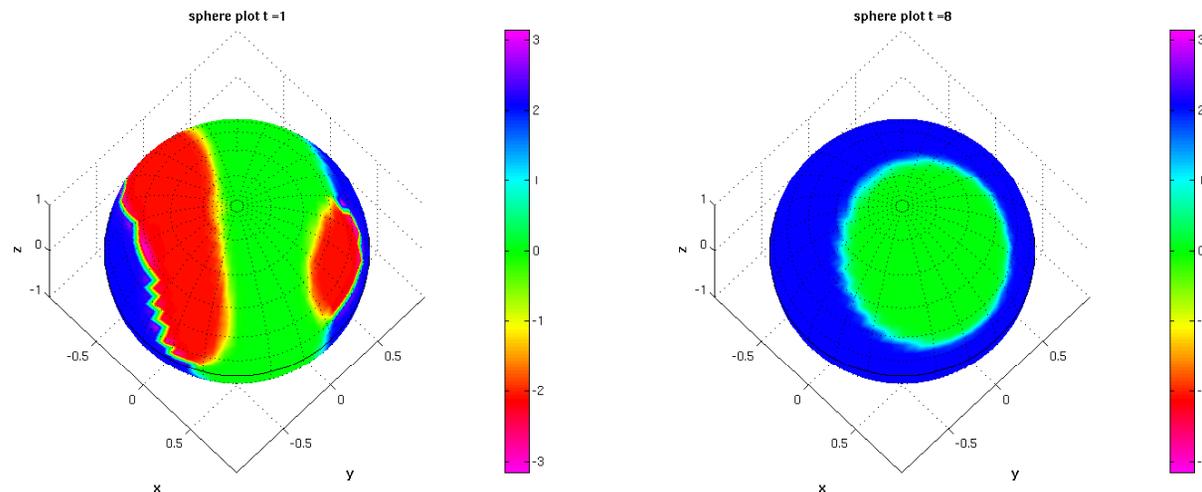
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# Application: Phase-Field Model



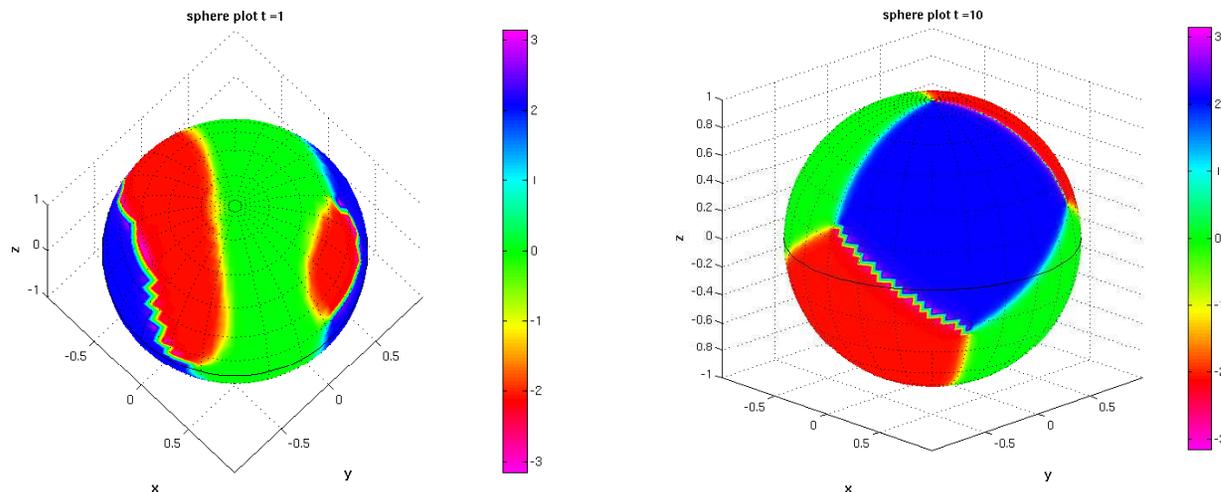
## Diffusion Fronts & Triple Point Dynamics

- ▷ complex-valued FFTs: *fftSc.m* & *ifftSc.m*
- ▷ complex-valued gradient flow,  $\psi(\lambda, \phi, t) \rightarrow$  diffusion & tri-stability

$$\psi_t = \frac{\delta \mathcal{F}}{\delta \psi^*} \quad ; \quad \mathcal{F}[\psi, \psi^*] = \int_S \left\{ \epsilon |\vec{\nabla} \psi|^2 + \frac{1}{\epsilon} |\psi - z_1|^2 |\psi - z_2|^2 |\psi - z_3|^2 \right\} dS$$

- ▷  $z_j^3 = 1$   $\rightarrow$  phases defined by 3 cube roots of unity
- ▷  $0 < \epsilon \ll 1$   $\rightarrow$  sharp fronts separating phases
- ▷ slow drift of phase fronts  $\rightarrow$  conformity
- ▷ on the plane, steady boundary-supported  $120^\circ$  triple junctions [Bronsard/Reitich, 1993]

# Application: Phase-Field Model



## Diffusion Fronts & Triple Point Dynamics

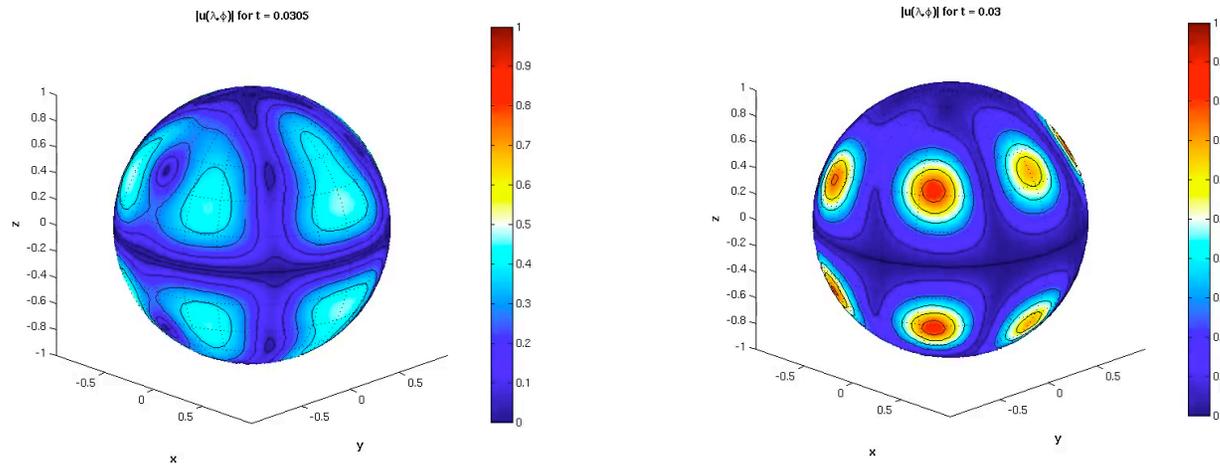
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# Application: Nonlinear Schrödinger Equation

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## Waves with Dispersion & Nonlinearity

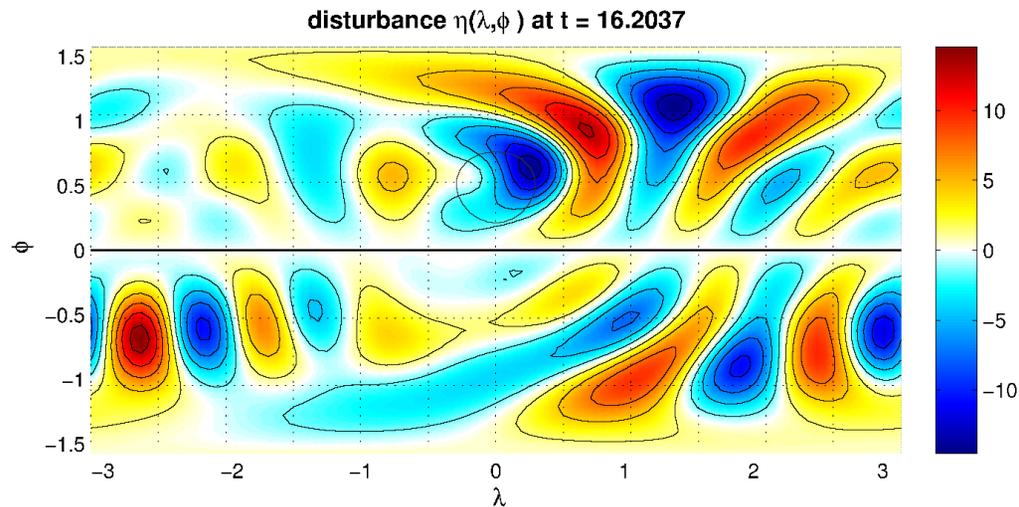
- ▷ complex-valued,  $\psi(\lambda, \phi, t)$

$$i\psi_t = \nabla^2 \psi \mp |\psi|^2 \psi$$

- ▷  $- \rightarrow$  defocussing;  $+ \rightarrow$  focussing
- ▷ near spherical harmonic initial condition
- ▷ on the plane: singularity possible in focussing case

## In Closing

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## Simple & Spectrally-Fast PDE Computing

- ▷ Fourier-based spectral transform
- ▷ implicit-explicit timestepping schemes
- ▷ suite of matlab routines
- ▷ follows paradigm for Fourier spectral method
- ▷ stability issues for fluid flows
- ▷ advection of rotating shallow water potential vorticity, DJM & Blazenko