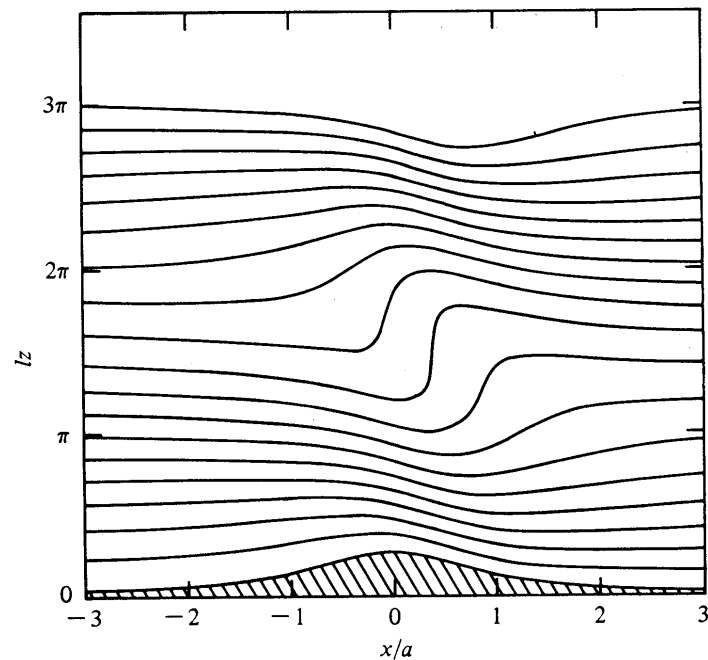
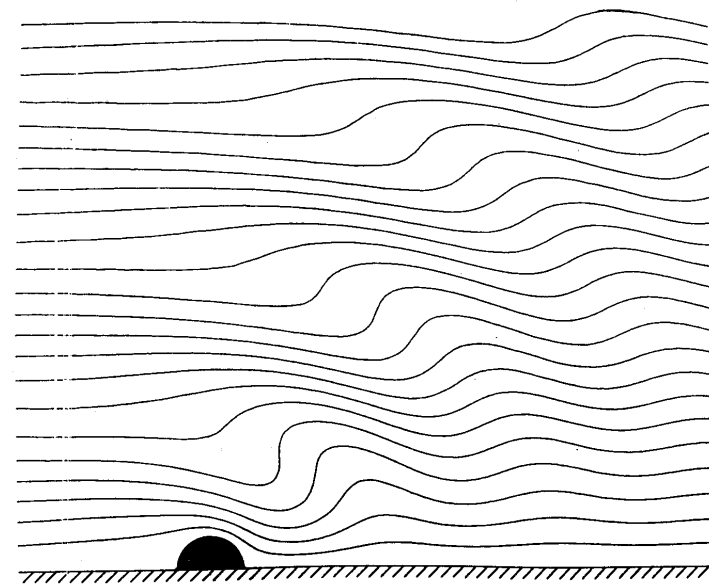


# A Few Surprises in 2D Nonlinear Flow over Topography

- ▷ computing Long's theory with exact surface boundary conditions



Lilly/Klemp 1979



Miles/Huppert 1968

- ▷ Dave Muraki (Simon Fraser University)

# Revisiting Long's 1953 Theory

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## Two-Dimensional Primitive Equations

- ▷ inviscid, incompressible, Boussinesq buoyancy

$$\begin{aligned}u_x + w_z &= 0 \\ \frac{Du}{Dt} &= -\phi_x \\ \delta^2 \frac{Dw}{Dt} - \theta &= -\phi_z \\ \frac{D\theta}{Dt} + w &= 0\end{aligned}$$

- ▷ nonhydrostatic parameter ( $\delta = U/NL$ ) & height scale ( $\mathcal{A} = NH/U$ )
- ▷ potential temperature ( $\theta$ ) & geopotential ( $\phi$ )

Steady Streamfunction:  $\psi(x, z) = z + \tilde{\psi}(x, z)$

- ▷ uniform upstream wind  $U$  & constant stratification  $N$
- ▷ exact reduction to linear Helmholtz equation for disturbance streamfunction

$$\delta^2 \tilde{\psi}_{xx} + \tilde{\psi}_{zz} + \tilde{\psi} = 0$$

- ▷ topographic surface at  $z = \mathcal{A}h(x)$  & streamline condition  $\rightarrow \psi(x, \mathcal{A}h(x)) = 0$
- ▷ radiation/decay BCs aloft

# Long 1955: Theory & Experiment

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$$\delta^2 \tilde{\psi}_{xx} + \tilde{\psi}_{zz} + \tilde{\psi} = 0$$

## Finite Amplitude Topography

- ▷ on streamline boundaries:  $\psi = Ah(x) + \tilde{\psi}(x, Ah(x)) = \text{constant}$

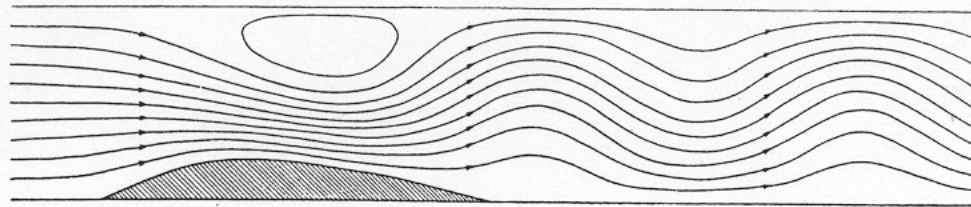
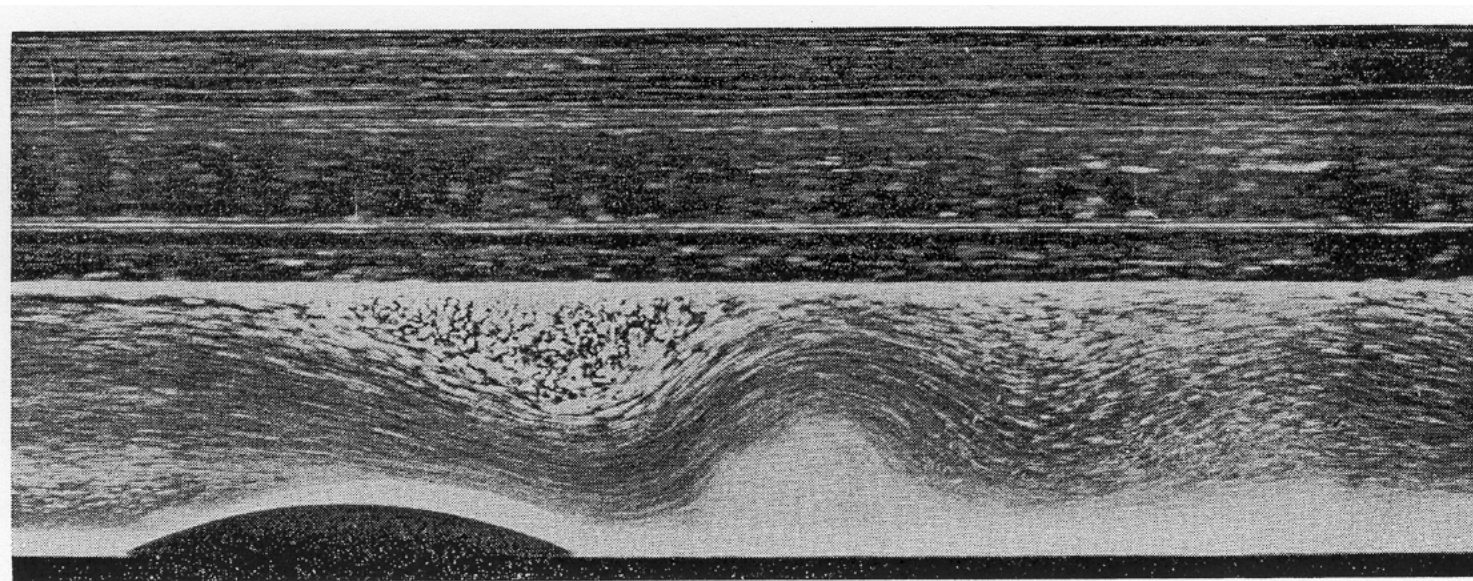


Fig. 8. Observed and calculated flow over an obstacle. Theoretical:  $F_1 = .200$ ,  $\delta = 1.0$ ,  $\alpha = .32$ . Experimental:  $F_1 = .204$ ,  $\delta = .200$ ,  $\alpha = .86$ .

Long 1953

# Linearized Surface Condition

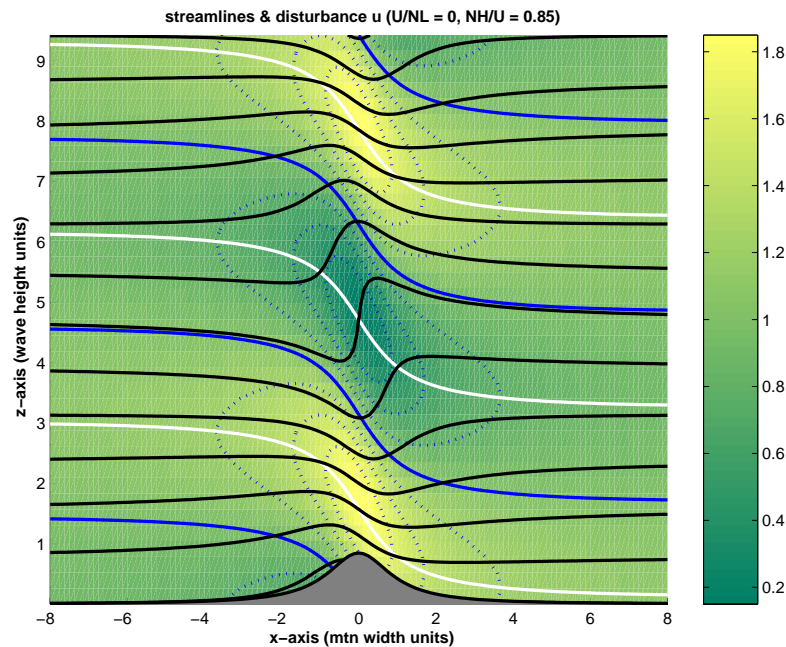
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$$\tilde{\psi}(x, z) = -\frac{\mathcal{A}}{2\pi} \int_{-\infty}^{+\infty} \hat{h}(k) e^{i(kx + m(k)z)} dk$$

## Fourier Solution (for small $\mathcal{A}$ )

- ▷ boundary at  $z = 0$  & linearized topographic condition  $\rightarrow \mathcal{A}h(x) + \tilde{\psi}(x, 0) = 0$
- ▷ aloft conditions via vertical mode number ( $\delta^2 k^2 + m^2 = 1$ )

$$m(k) = \begin{cases} \text{sign}(k) \sqrt{1 - \delta^2 k^2} & \text{for } |\delta k| \leq 1 \text{ (long scale radiation)} \\ i \sqrt{\delta^2 k^2 - 1} & \text{for } |\delta k| \geq 1 \text{ (short scale decay)} \end{cases}$$



## General Helmholtz Solution

---

$$\tilde{\psi}(x, z) = -\mathcal{A} \int_{-\infty}^{+\infty} \hat{c}(k) e^{i(kx + m(k)z)} dk$$

### Fourier Representation

- ▷ satisfies aloft conditions ( $\delta^2 k^2 + m^2 = 1$ )
- ▷ surface at  $z = \mathcal{A}h(x)$  & exact topographic condition  $\rightarrow \mathcal{A}h(x) + \tilde{\psi}(x, \mathcal{A}h(x)) = 0$

$$h(x) - \int_{-\infty}^{+\infty} \hat{c}(k) e^{i(kx + m(k)\mathcal{A}h(x))} dk = 0$$

### Fredholm Integral Equation of the First-Kind

- ▷ linearity: action of integral operator is linear in unknown coefficients  $\hat{c}(k)$
- ▷ numerical solution equivalent to matrix inversion
- ▷ velocities from spectral differentiation:  $u = \psi_z$  &  $w = -\psi_x$
- ▷ no need to compute Fourier transform:  $c(x) \rightarrow$  *effective topography*

## Direct Steady Solve

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$$h(x) - \int_{-\infty}^{+\infty} \hat{c}(k) e^{i(kx + m(k)Ah(x))} dk = 0$$

### Numerical Discretization

- ▷ collocation points:  $\{x_1 \dots x_\alpha \dots x_N\}$  &  $N$  knowns:  $h_\alpha = h(x_\alpha)$
- ▷ wavenumbers:  $\{k_1 \dots k_\beta \dots k_N\}$  &  $N$  unknowns:  $\hat{c}_\beta \approx \hat{c}(k_\beta)$
- ▷ approximate integral at each  $x_\alpha$  by quadrature (trapezoidal rule) over  $\beta = 1 \dots N$

$$h_\alpha - \sum_{\beta=1}^N \hat{c}_\beta \underbrace{e^{i(k_\beta x_\alpha + m(k_\beta)Ah(x_\alpha))}}_{\mathbf{K}_{\alpha,\beta}} w_\beta \Delta k = 0$$

### Matrix Inversion

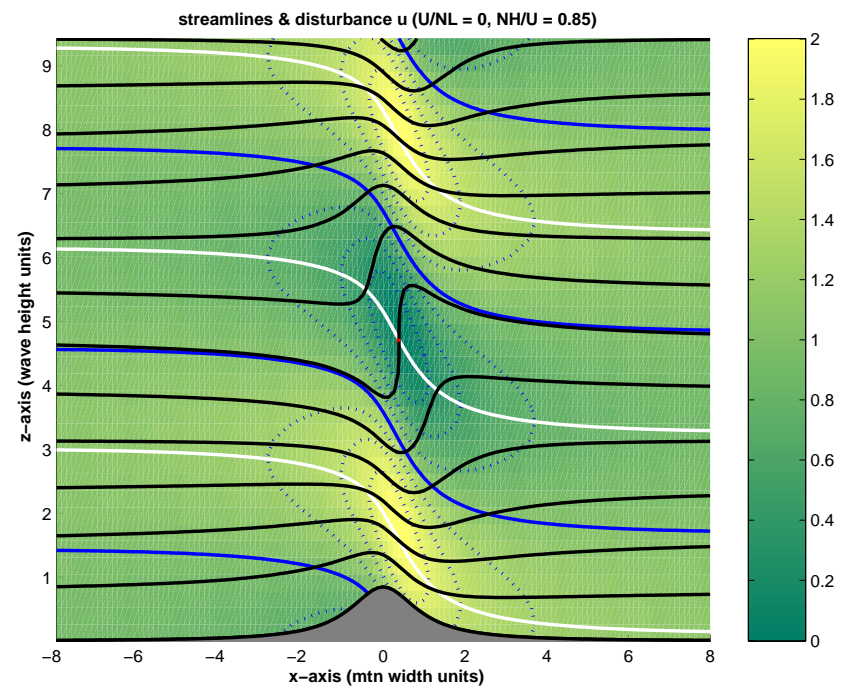
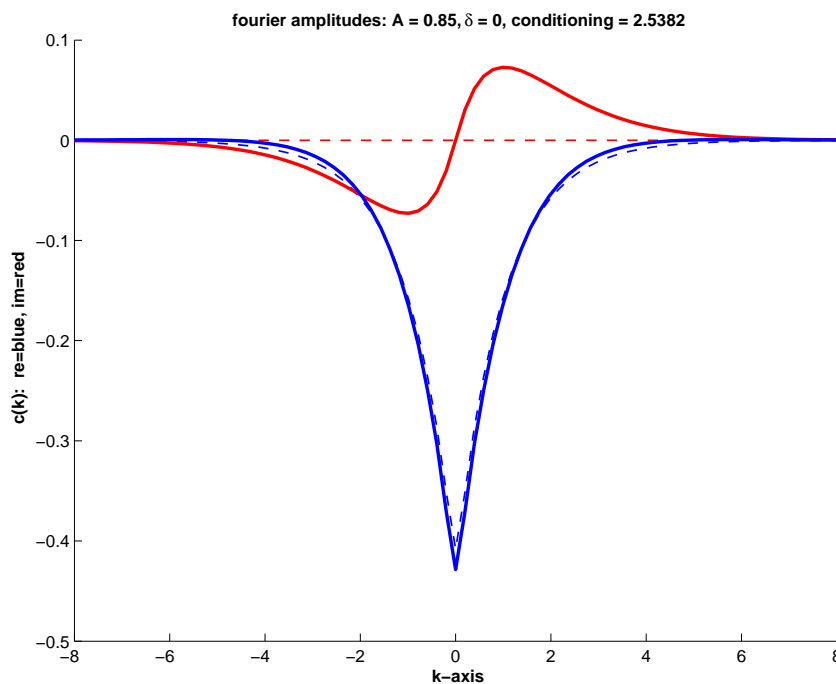
- ▷  $N$  linear equations in  $N$  unknowns:  $(\vec{h}_\alpha) = [\mathbf{K}_{\alpha,\beta}] (\vec{c}_\beta)$
- ▷  $m(k)$  is discontinuous at  $k = 0 \rightarrow$  half-line integrals
- ▷ full matrix  $\mathbf{K}$  can be ill-conditioned  $\rightarrow$  catastrophic loss of precision as  $N$  increases



# Numerical Implementation

## Fourier Conditioning

- ▷  $\mathcal{A} = 0$  recovers linear theory & discrete Fourier transform is well-conditioned
- ▷ equi-spaced discretizations with  $\Delta k \Delta x = 2\pi/N$  is essential

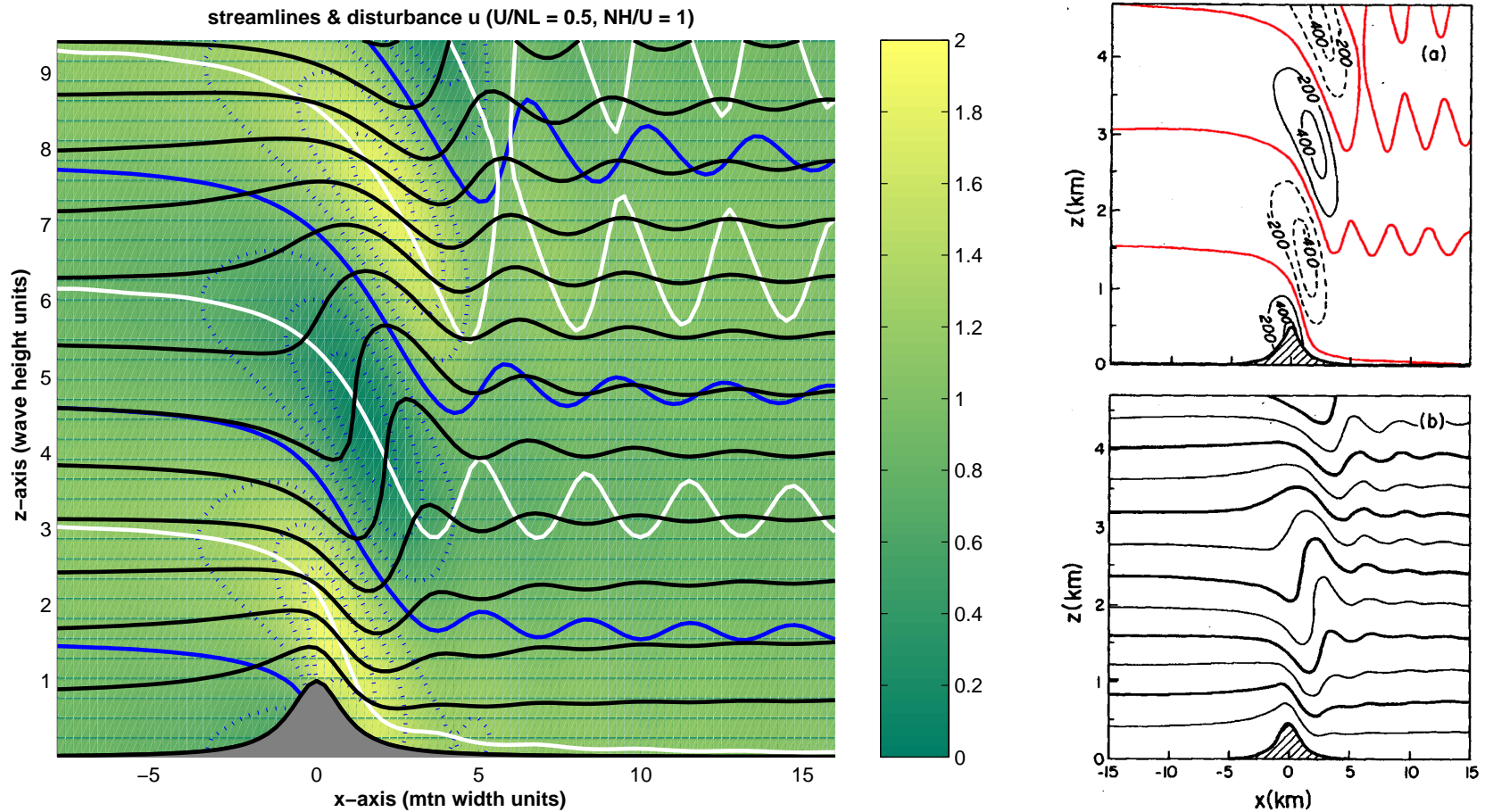


- ▷ Lilly/Klemp 1979, hydrostatic critical overturning ( $\mathcal{A} \approx 0.85$ )  
→  $N = 256$ , 1.1s to solve & 2.0s to plot
- ▷ Fourier representation allows periodic wrap-around → large computational domains

# A Nonhydrostatic Example

Laprise & Peltier, 1988

- ▷ predictor/corrector to obtain effective topography  $c(x)$  → typically 50 iterations



- ▷ large amplitude  $\mathcal{A} = 1.0$  & moderately nonhydrostatic  $\delta = 0.5$
- ▷  $N = 2056$ ,  $x_\infty = 256$ : 284s to solve, 89s to plot, log-condition number = 5.75

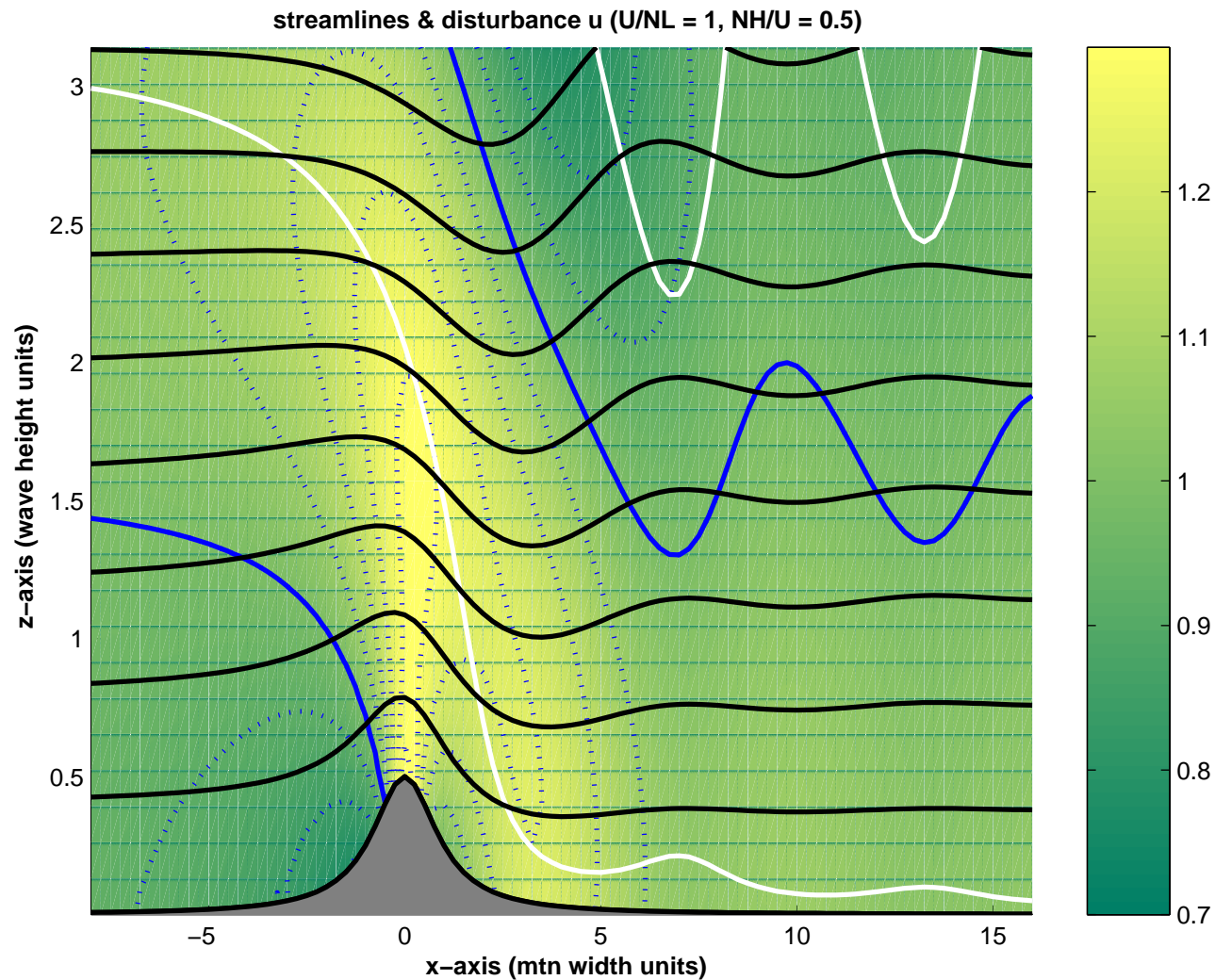


# A Strongly Nonhydrostatic Example

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$$\delta = 1.0 \quad \& \quad \mathcal{A} = 0.5$$

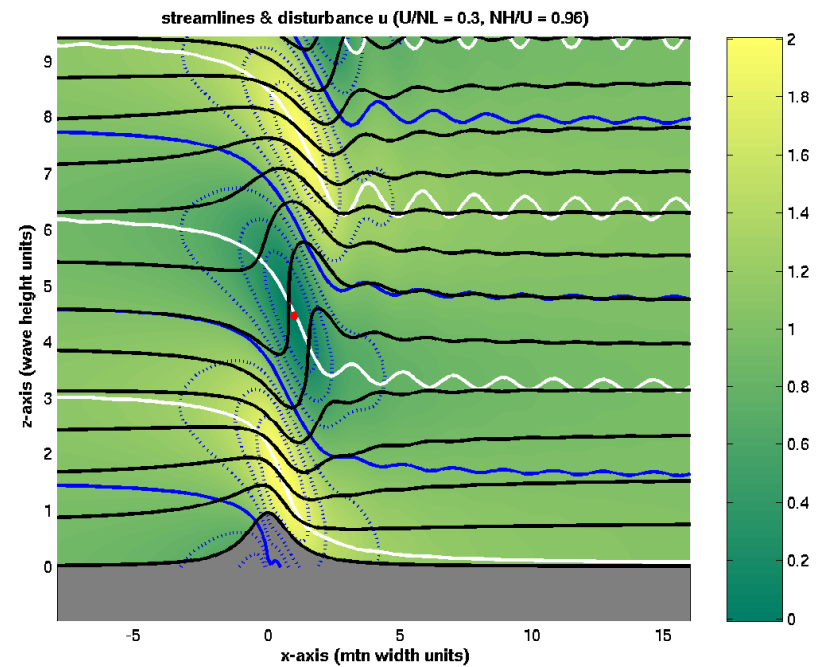
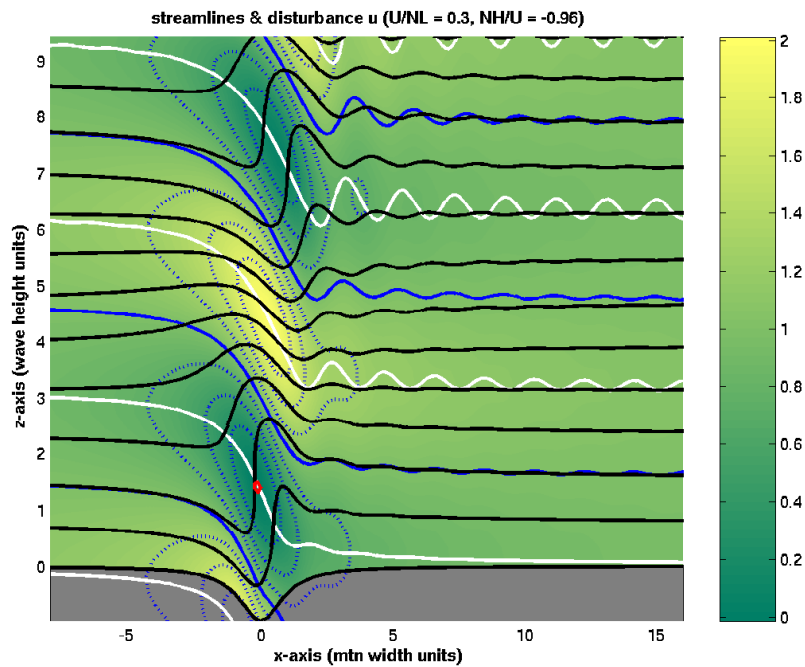
- ▷  $u$ -wind maximum shifts towards the summit as nonhydrostatic effect increases



# Mountain vs Valley

$$\delta = 0.3 \quad \& \quad \mathcal{A} = \pm 0.96$$

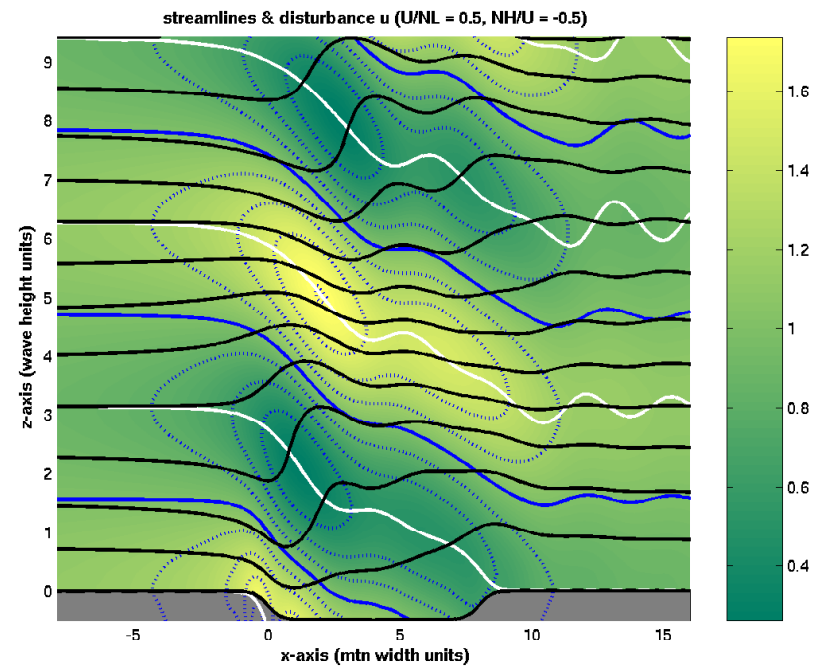
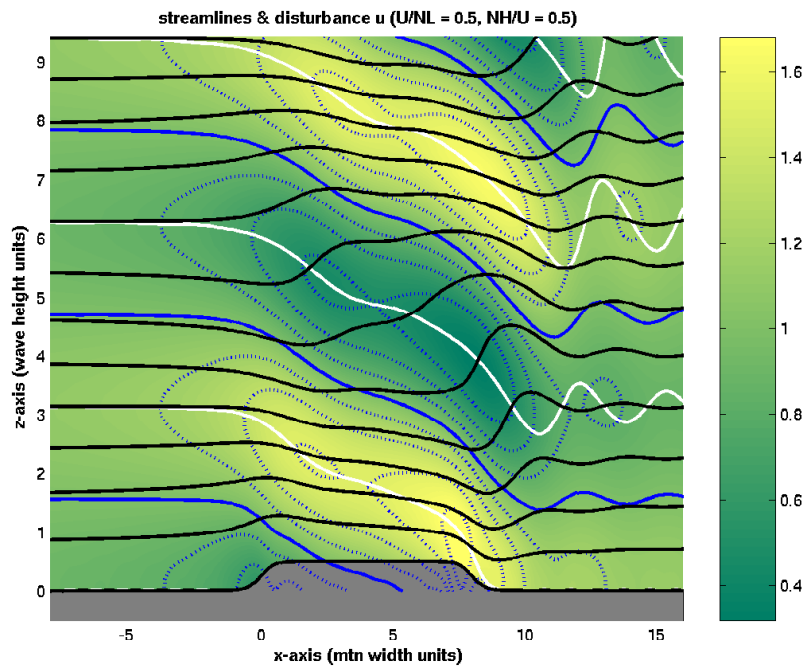
- ▷ both cases near critical overturning
- ▷ very little asymmetry in overall magnitude of response (unlike rotating case)



# Extended Topography

$$\delta = 0.5 \quad \& \quad \mathcal{A} = \pm 0.5$$

- ▷ largest response associated with downslope
- ▷ slightly more wind in valley case:  $0.32 < u^+ < 1.77$  vs  $0.26 < u^- < 1.82$



# Lyra's Topographic Greens Function

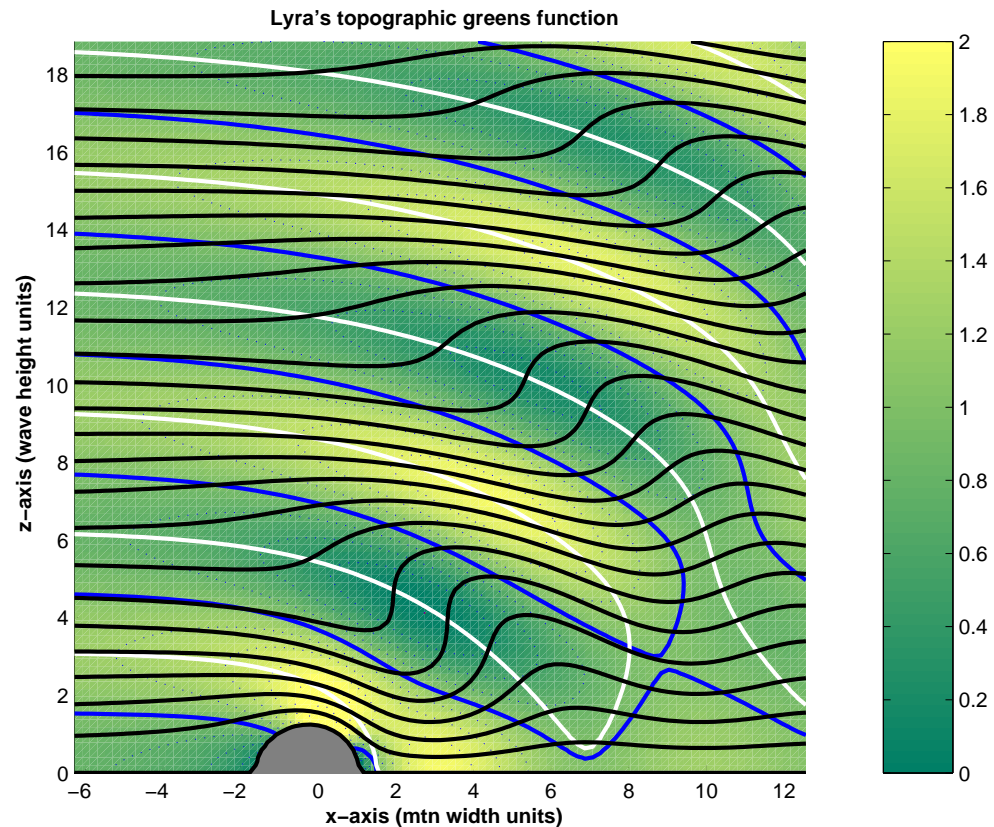
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## Delta Function Topography

- ▷ from Lyra 1940 & 1943 (via Alaka 1960) for  $\delta = 1$ : Bessel series

$$\tilde{\psi}(r, \theta) = \frac{1}{2} Y_1(r) \sin \theta + \frac{1}{\pi} \sum_1^{\infty} \frac{4n}{4n^2 - 1} J_{2n}(r) \sin 2n\theta$$

- ▷ critical overturning for delta strength  $\approx 4.06$

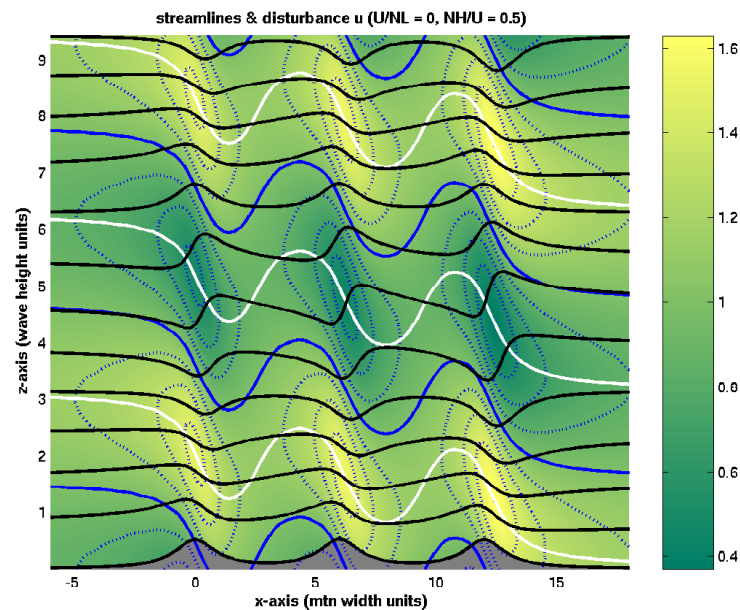


# Topographic Boundary Conditions

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## Direct Computation

- ▷ consistent with spectral radiation condition
- ▷ elementary formulation for non-iterative solve
- ▷ Long's theory: hydrostatic & nonhydrostatic
- ▷ Fredholm first-kind integral equation is generally ill-conditioned  
→ possible resolution via Lyra's greens function
- ▷ Fourier representation allows for wrap-around of waves



- ▷ open issues in stability of Long's solutions?