

OVERLAPPING SCHWARZ ITERATION

1. INTRODUCTION

Our goal is to either analytically solve, or numerically approximate, the 2nd Neumann eigenfunction of the Laplacian on a generic acute triangle $\Omega \equiv ABC$ (see figure). Since we don't have an analytic solution on this domain, but can solve problems on sectors or the circle, we propose a domain decomposition strategy in the spirit of the overlapping Schwarz iteration.

Let O be the incenter of the triangle, and denote by Ω_0 the interior of the incircle. This is tangent to the segments AB, BC and CA at F, D, E respectively. Denote by Γ_{10} the segment of the circle connecting E, F . Denote by Ω_1 the sector AEF , and by Γ_{01} the arc connecting E, F . Note that Ω_1 and Ω_2 overlap. Repeat this process for the other three vertices. Let us denote the exact second Neumann eigenvalue as λ .

Now we'll proceed by iteration. At step n , suppose u_i^n for $i = 1, 2, 3$ satisfies

$$\begin{aligned}
 (1) \quad & -\Delta u_i^n = \lambda_i^n u_i, \quad x \in \Omega_i \\
 (2) \quad & \frac{\partial u_i^n}{\partial \nu} = 0 \quad x \in \partial\Omega_i \setminus \Gamma_{0i} \\
 (3) \quad & u_0^{n-1} \frac{\partial u_i^n}{\partial \nu} - \frac{\partial u_0^{n-1}}{\partial \nu} u_i^n = 0, \quad x \in \Gamma_{0i} \\
 (4) \quad &
 \end{aligned}$$

The function u_0^n solves

$$\begin{aligned}
 (5) \quad & -\Delta u_0^n = \lambda_0^n u_0^n, \quad x \in \Omega_0 \\
 (6) \quad & u_i^n \frac{\partial u_0^n}{\partial \nu} - \frac{\partial u_i^n}{\partial \nu} u_0^n = 0, \quad x \in \Gamma_{i0} \\
 (7) \quad &
 \end{aligned}$$

Now (u_i^n, λ_i) solve generalized eigenfunction problems on sectors of circles. If one knows $u_0^{n-1}, \frac{\partial u_0^{n-1}}{\partial \nu}$ on the curves Γ_{0i} , then one can use local Fourier-Bessel expansions to get u_i^n . Then one uses their traces onto Γ_{10} , and has the correct data to solve the eigenvalue problem for u_0^n on the disk. One can use Fourier-Bessel expansions to do this as well.

The claim is that as $n \rightarrow \infty$, the sequences u_i^n converge to the restriction of the actual eigenfunction u on the sub-domains. Clearly we have to prescribe a starting guess for the iteration.

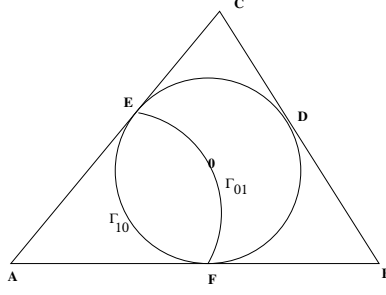


FIGURE 1. Domains

1.1. **Solving on the wedge Ω_i .** We are at step n of the iteration. Suppose $i = 1$ is fixed for concreteness, and the traces of u_0^{n-1} and $\frac{\partial u_0^{n-1}}{\partial \nu}$ on Γ_{01} are known.

$$(8) \quad -\Delta u_1^n = \lambda_1^n u_1, \quad x \in \Omega_1$$

$$(9) \quad \frac{\partial u_1^n}{\partial \nu} = 0 \quad x \in \partial\Omega_1 \setminus \Gamma_{01}$$

$$(10) \quad u_0^{n-1} \frac{\partial u_1^n}{\partial \nu} - \frac{\partial u_0^{n-1}}{\partial \nu} u_1^n = 0, \quad x \in \Gamma_{01}$$

(11)

We will try the method of particular solutions of Fox and Henrici, adapted to this problem. We know that since the opening angle is α , in Ω_1 the functions

$$(12) \quad w_k(r, \theta) := J_{\frac{\pi k}{\alpha}}(\sqrt{\lambda}r) \cos\left(\frac{\pi k}{\alpha}\theta\right)$$

will satisfy the Neumann conditions on the line segments AE, AF , as well as satisfy the equation $-\Delta w_k = \lambda w_k$.

So, we suppose $u_1^n = \sum_{k=1}^M c_k w_k(r, \theta)$. We want to find the coefficients c_k so as to satisfy the boundary condition on Γ_{01} . Now, Γ_{01} is an arc of radius $\rho_1 = AE$. Let (ρ_1, θ_j) be $2M$ collocation points along this curve. At each point, we want to enforce

$$0 \quad (\#3) u_0^{n-1}(\rho_1, \theta_j) \frac{\partial u_1^n}{\partial \nu}(\rho_1, \theta_j) - \frac{\partial u_0^{n-1}}{\partial \nu}(\rho_1, \theta_j) u_1^n(\rho_1, \theta_j)$$

$$(\#4) u_0^{n-1}(\rho_1, \theta_j) \sum_{k=1}^M c_k \frac{\partial}{\partial r} J_{\frac{\pi k}{\alpha}}(\sqrt{\lambda}\rho_1) \cos\left(\frac{\pi k}{\alpha}\theta_j\right) - \frac{\partial u_0^{n-1}}{\partial \nu}(\rho_1, \theta_j) \sum_{k=1}^M c_k J_{\frac{\pi k}{\alpha}}(\sqrt{\lambda}\rho_1) \cos\left(\frac{\pi k}{\alpha}\theta_j\right)$$

This is equivalent to solving the rectangular nonlinear system

$$(15) \quad A(\lambda)\vec{c} = 0$$

where $a_{jk}(\lambda) = u_0^{n-1}(\rho_1, \theta_j) \frac{\partial}{\partial r} J_{\frac{\pi k}{\alpha}}(\sqrt{\lambda}\rho_1) \cos\left(\frac{\pi k}{\alpha}\theta_j\right) - \frac{\partial u_0^{n-1}}{\partial \nu}(\rho_1, \theta_j) J_{\frac{\pi k}{\alpha}}(\sqrt{\lambda}\rho_1) \cos\left(\frac{\pi k}{\alpha}\theta_j\right)$

We find the solutions by looking for values of λ so that the smallest singular value of $A(\lambda)$ approaches 0. This is the Moler approach to the original Fox-Henrici-Moler paper.

Once we locate the solution \vec{c} , we have the iterate u_1^n . We do this same process for the other wedges as well.

1.2. Solving on the disk Ω_0 . We are at step n of the iteration, and have solved for the functions u_i^n . We therefore have their traces on the arcs Γ_{i0} .

The function u_0^n solves

$$(16) \quad -\Delta u_0^n = \lambda_0^n u_0^n, \quad x \in \Omega_0$$

$$(17) \quad u_i^n \frac{\partial u_0^n}{\partial \nu} - \frac{\partial u_i^n}{\partial \nu} u_0^n = 0, \quad x \in \Gamma_{i0}$$

(18)

We shall again use a Fourier-Bessel ansatz: let $z_m(r, \theta) = J_m(\lambda r)e^{im\theta}$, and assume

$$u_0^n(r, \theta) = \sum_{k=0}^M d_k z_k(r, \theta)$$

. Repeat the process above of enforcing the (non-standard) boundary conditions at collocation points along Γ_{i0} .

2. CONVERGENCE?

At the end of the n th step, we have 4 functions: u_i^n , $i=0,1,2,3$ and 4 eigenvalues. We repeat the process until the eigenvalues are all the same number.