

## Vector Functions of a Single Variable

Definition: Vector-valued function  $\vec{F}(t)$  :  
 assigns a vector  $\vec{F}$  to each  $t$  in some interval  $t_1 \leq t \leq t_2$

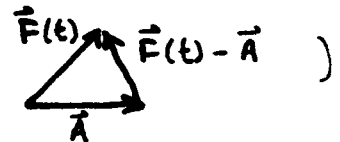
$$\vec{F}(t) = F_1(t) \hat{i} + F_2(t) \hat{j} + F_3(t) \hat{k}$$

Limit:  $\lim_{t \rightarrow t_0} \vec{F}(t) = \vec{A}$

$$\forall \epsilon > 0 \quad \exists \delta > 0$$

$$\text{s.t. } 0 < |t - t_0| < \delta \Rightarrow |\vec{F}(t) - \vec{A}| < \epsilon$$

(ie  $|\vec{F}(t)| \rightarrow |\vec{A}|$ , and angle between  $\vec{F}(t)$  and  $\vec{A}$  approaches zero.



Equivalent:

Components of  $\vec{F}(t)$  approach components of  $\vec{A}$   $\lim_{t \rightarrow t_0} F_i(t) = A_i$   $i=1,2,3$

Continuity:  $\vec{F}(t)$  is continuous at  $t_0$  if

$$\lim_{t \rightarrow t_0} \vec{F}(t) = \vec{F}(t_0) \Leftrightarrow \lim_{t \rightarrow t_0} F_i(t) = F_i(t_0)$$

Derivative:  $\vec{F}(t)$  is differentiable at  $t_0$  if the limit

$$\lim_{\Delta t \rightarrow 0} \frac{\vec{F}(t_0 + \Delta t) - \vec{F}(t_0)}{\Delta t} = \vec{F}'(t_0) = \frac{d\vec{F}}{dt}(t_0)$$

exists

$$\Leftrightarrow F_i'(t) \text{ exists, } i=1,2,3$$

Note: for a vector written in terms of its components w.r.t a basis:

Differentiate componentwise:

$$\vec{F}'(t) = F_1'(t) \hat{i} + F_2'(t) \hat{j} + F_3'(t) \hat{k}$$

$$\left( \frac{d\vec{F}}{dt} \right)_i = \frac{dF_i}{dt}$$

## Differentiation Rules:

Theorem: If  $\vec{F}$  and  $\vec{G}$  are differentiable vector functions and  $s$  is a differentiable scalar function, then

sum rule 1.  $\frac{d}{dt} (\vec{F} + \vec{G}) = \frac{d\vec{F}}{dt} + \frac{d\vec{G}}{dt}$

product rules  $\left\{ \begin{array}{l} 2. \frac{d}{dt} (s\vec{F}) = \frac{ds}{dt} \vec{F} + s \frac{d\vec{F}}{dt} \end{array} \right.$

3.  $\frac{d}{dt} (\vec{F} \cdot \vec{G}) = \frac{d\vec{F}}{dt} \cdot \vec{G} + \vec{F} \cdot \frac{d\vec{G}}{dt}$

4.  $\frac{d}{dt} (\vec{F} \times \vec{G}) = \frac{d\vec{F}}{dt} \times \vec{G} + \vec{F} \times \frac{d\vec{G}}{dt}$

← note order!

chain rule 5.  $\frac{d}{dt} (\vec{F}(s(t))) = s'(t) \vec{F}'(s(t))$

Proofs: eg 4.  $\frac{d}{dt} (F_i G_i) = \frac{dF_i}{dt} G_i + F_i \frac{dG_i}{dt}$

5.  $\frac{d}{dt} (\vec{F} \times \vec{G})_i = \frac{d}{dt} (\epsilon_{ijk} F_j G_k) = \epsilon_{ijk} \left( \frac{dF_j}{dt} G_k + F_j \frac{dG_k}{dt} \right)$   
 $= \epsilon_{ijk} \frac{dF_j}{dt} G_k + \epsilon_{ijk} F_j \frac{dG_k}{dt} = \left( \frac{d\vec{F}}{dt} \times \vec{G} \right)_i + \left( \vec{F} \times \frac{d\vec{G}}{dt} \right)_i$

eg Suppose  $\vec{F}(t)$  has constant magnitude (varies only in direction)

$|\vec{F}(t)| = c$  i.e.  $\vec{F}(t)$  traces out a curve on the surface of a sphere of radius  $c$ .

Then either  $\vec{F}'(t) = \vec{0}$  or  $\vec{F}'(t) \perp \vec{F}(t)$

Proof.  $|\vec{F}(t)| = \text{const} = c \Rightarrow |\vec{F}(t)|^2 = c^2$   
 $\Rightarrow \vec{F}(t) \cdot \vec{F}(t) = c^2$

$\Rightarrow \frac{d}{dt} |\vec{F}(t)|^2 = \frac{d}{dt} \vec{F}(t) \cdot \vec{F}(t) = 0$

$\Rightarrow \vec{F}(t) \cdot \frac{d\vec{F}}{dt} + \frac{d\vec{F}}{dt} \cdot \vec{F}(t) = 0 \Rightarrow 2 \vec{F} \cdot \frac{d\vec{F}}{dt} = 0$

i.e. if  $\frac{d\vec{F}}{dt} \neq \vec{0}$ , then  $\frac{d\vec{F}}{dt} \perp \vec{F}$  (tangent to curve is  $\perp$  radius of sphere)

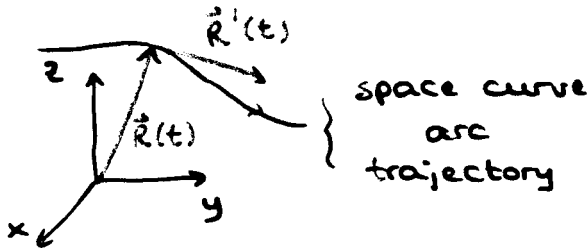
# Space Curves

$$\vec{R} = \vec{R}(t) = x(t)\hat{i} + y(t)\hat{j} + z(t)\hat{k}$$

$t_1 \leq t \leq t_2$   $\vec{R}(t)$  continuous

$\vec{R}(t)$  is the position vector of a point in space

As  $t$  varies,  $\vec{R}(t)$  traces out a space curve



a one-dimensional object:  
a single value  $t$  is needed to specify each point on the arc

$t$ : parameter

$\vec{R}(t)$ : parametrization of the curve

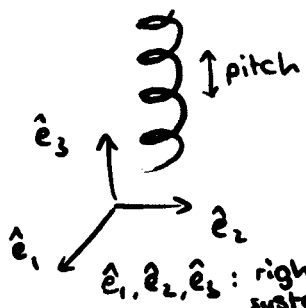
eg.  $\vec{R}(t) = \vec{R}_0 + t\vec{V}$ ,  $-\infty < t < \infty$  line  
 $\vec{R}_0$ : fixed point on line  
 $\vec{V}$ : vector parallel to line

$\vec{R}(t) = \vec{R}_0 + (\tan t)\vec{W}$ ,  $-\pi/2 < t < \pi/2$ ,  $\vec{W} = \frac{1}{3}\vec{V}$   
 alternative parametrization of line

$\vec{R}(t) = \cos t \hat{i} + \sin t \hat{j}$  circle radius 1, in x-y plane

$\vec{R}(t) = \vec{R}_0 + p \cos t \hat{i} + p \sin t \hat{j}$  circle center  $\vec{R}_0$ , radius  $p$   
 (in plane  $\perp \hat{k}$  i.e.  $\parallel$  to x-y plane)  
 - a parametrization of  $(x-x_0)^2 + (y-y_0)^2 = p^2$ ,  $z = z_0$

$\vec{R}(t) = \cos t \hat{i} + \sin t \hat{j} + t \hat{k}$  helix projection onto x-y plane is unit circle



$\vec{R}(t) = \vec{R}_0 + p \cos t \hat{e}_1 + p \sin t \hat{e}_2 + at \hat{e}_3$   
 helix, radius  $p$ , with axis through  $\vec{R}_0$ , parallel to  $\hat{e}_3$   
 $2\pi/|a|$ , right-handed if  $\{a > 0$   
 left-handed if  $\{a < 0$

$\hat{e}_1, \hat{e}_2, \hat{e}_3$ : right-handed system of unit vectors

Common interpretation:  $t$ : time,  $\vec{R}(t)$ : trajectory of particle position vector

In time interval  $(t, t + \Delta t)$ :

• Displacement  $\Delta \vec{R} = \vec{R}(t + \Delta t) - \vec{R}(t) = \Delta x \hat{i} + \Delta y \hat{j} + \Delta z \hat{k}$

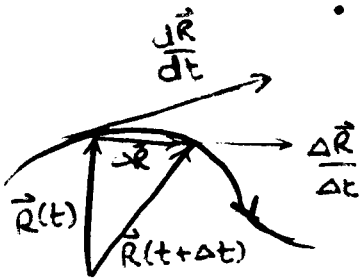
$\Rightarrow \frac{\Delta \vec{R}}{\Delta t} = \frac{\Delta x}{\Delta t} \hat{i} + \frac{\Delta y}{\Delta t} \hat{j} + \frac{\Delta z}{\Delta t} \hat{k}$  average velocity

(Instantaneous)  
• Velocity  $\vec{v}(t) = \vec{R}'(t) = \frac{d\vec{R}}{dt} = \lim_{\Delta t \rightarrow 0} \frac{\Delta \vec{R}}{\Delta t}$

$\Rightarrow \vec{v}(t) = \frac{dx}{dt} \hat{i} + \frac{dy}{dt} \hat{j} + \frac{dz}{dt} \hat{k}$

• Acceleration  $\vec{a}(t) = \frac{d^2 \vec{R}}{dt^2} = \vec{R}''(t) = \frac{d^2 x}{dt^2} \hat{i} + \frac{d^2 y}{dt^2} \hat{j} + \frac{d^2 z}{dt^2} \hat{k}$

• Speed  $v(t) = |\vec{v}(t)| = \left[ \left( \frac{dx}{dt} \right)^2 + \left( \frac{dy}{dt} \right)^2 + \left( \frac{dz}{dt} \right)^2 \right]^{1/2}$   
 $(= \frac{ds}{dt}) = \left( \frac{dR_i}{dt} \frac{dR_i}{dt} \right)^{1/2}$



Note:  $\vec{v}(t_0) = \frac{d\vec{R}}{dt}(t_0)$  is tangent to the curve  
 $\vec{R} = \vec{R}(t)$  at  $t = t_0$ .

Unit tangent  $\vec{T}(t) = \frac{\vec{R}'(t)}{|\vec{R}'(t)|} = \frac{\vec{v}(t)}{v(t)}$

$\vec{T} = \frac{\frac{dx}{dt} \hat{i} + \frac{dy}{dt} \hat{j} + \frac{dz}{dt} \hat{k}}{\sqrt{\left( \frac{dx}{dt} \right)^2 + \left( \frac{dy}{dt} \right)^2 + \left( \frac{dz}{dt} \right)^2}}$

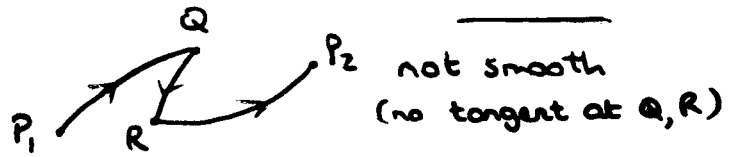
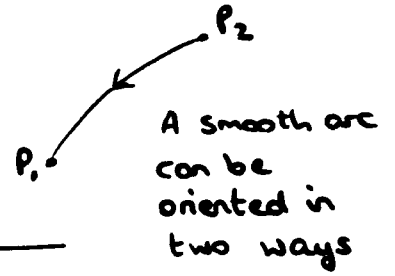
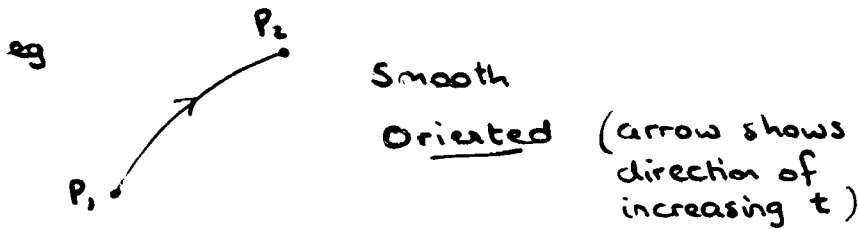
Defn: An arc is smooth if it has a parametrization  $\vec{R} = \vec{R}(t)$ ,  
 $t_1 \leq t \leq t_2$ , satisfying: [possibly  $t_1 = -\infty, t_2 = \infty$ ]

i)  $\vec{R}'(t)$  exists and is continuous for  $t_1 \leq t \leq t_2$

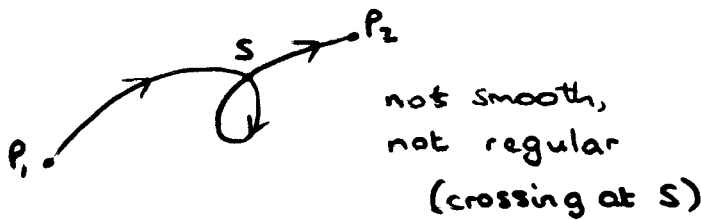
ii) No self-intersections: distinct  $t, t_1 < t < t_2$ , correspond to different points [ie  $\vec{R}(t_1) = \vec{R}(t_2) \Rightarrow \theta_1 = \theta_2$   
 $t_1 < \theta_1, \theta_2 < t_2$ ]

iii)  $\vec{R}'(t) \neq \vec{0}$ ,  $t_1 \leq t \leq t_2$

[ie there is a tangent at every point: no cusps]



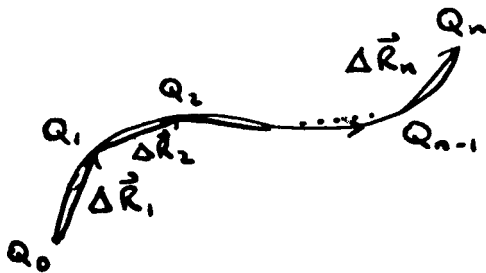
regular (union of finite no. of smooth arcs, no self-intersections)



Note: The straight line  $\vec{R}(t) = \vec{R}_0 + t^3 \vec{v}$ ,  $-1 \leq t \leq 1$  violates condition iii), since  $\vec{R}'(t) = \vec{0}$  at  $t=0$  but it is a smooth curve since it can be parametrized by  $\vec{R}(t) = \vec{R}_0 + t \vec{v}$ ,  $-1 \leq t \leq 1$

Arc Length

$C: \vec{R} = \vec{R}(t)$ ,  $a \leq t \leq b$  smooth space curve



Partition  $[a, b]$  into  $n$  subintervals

$$a = t_0 < t_1 < t_2 < \dots < t_{n-1} < t_n = b$$

$Q_j$  has position vector  $\vec{R}(t_j)$ ,  $0 \leq j \leq n$

Length of each segment

$$\begin{aligned} |\Delta \vec{R}_j| &= |\vec{R}(t_j) - \vec{R}(t_{j-1})| = |\Delta x_j \hat{i} + \Delta y_j \hat{j} + \Delta z_j \hat{k}| \\ &= (\Delta x_j^2 + \Delta y_j^2 + \Delta z_j^2)^{1/2} \end{aligned}$$

Mean value theorem:  $\Delta x_j = x(t_j) - x(t_{j-1}) = \frac{dx}{dt}(t_j) \Delta t_j$   
 for some  $\tau_j$ ,  $t_{j-1} \leq \tau_j \leq t_j$   $\Delta t_j = t_j - t_{j-1}$

and  $\Delta y_j = \frac{dy}{dt}(\tau_j') \Delta t_j$ ,  $\Delta z_j = \frac{dz}{dt}(\tau_j'') \Delta t_j$

Length of polygonal path

$$\sum_{j=1}^n |\Delta \vec{R}_j| = \sum_{j=1}^n \left[ \left( \frac{dx}{dt}(\tau_j) \right)^2 + \left( \frac{dy}{dt}(\tau_j') \right)^2 + \left( \frac{dz}{dt}(\tau_j'') \right)^2 \right]^{1/2} \Delta t_j$$

Take limit as maximum length of subdivision  $\rightarrow 0$ :  $\max_{1 \leq j \leq n} |\Delta t_j| \rightarrow 0$

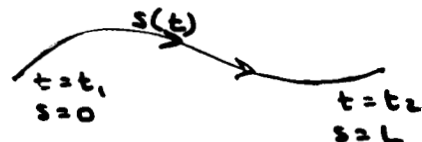
Riemann integral:

$\Rightarrow$  Arc Length  $L = \int_a^b \left[ \left( \frac{dx}{dt} \right)^2 + \left( \frac{dy}{dt} \right)^2 + \left( \frac{dz}{dt} \right)^2 \right]^{1/2} dt$   
 $= \int_a^b \left| \frac{d\vec{R}}{dt} \right| dt$

$$L = \int_C |d\vec{R}| = \int_C [dx^2 + dy^2 + dz^2]^{1/2}$$

$\nearrow$  arc length is a property of the curve, independent of parametrization  $\approx dx = \frac{dx}{dt} dt$   
...

Arc length measured along curve



$$s = s(t) = \int_{t_1}^t \left| \frac{d\vec{R}}{dt} \right| dt$$

$$t \geq t_1$$

Note:  $\frac{ds}{dt} = \left| \frac{d\vec{R}}{dt} \right| = |\vec{v}(t)| = v(t)$  : Speed = rate of change of arclength

$$" ds = |d\vec{R}| = (dx^2 + dy^2 + dz^2)^{1/2} "$$

C a smooth curve  $\Rightarrow \vec{R}'(t) \neq \vec{0}$  (for some parametrization)

$$\Rightarrow \frac{ds}{dt} = v(t) = |\vec{R}'(t)| > 0 \quad (\neq 0)$$

ie the function  $s = s(t)$  is monotonic, one-to-one

$\Rightarrow s(t)$  can (in principle) be inverted to give  $t = t(s)$

$\Rightarrow$  Reparametrize - use  $s$  as a parameter, substitute  $t(s)$

Arc length parametrization:  $\vec{R}(s)$  (in principle; direct computation usually difficult)

Note:  $\frac{d\vec{R}}{dt} = \frac{d\vec{R}}{ds} \frac{ds}{dt} \Rightarrow \frac{d\vec{R}}{ds} = \frac{d\vec{R}}{dt} / \frac{ds}{dt} = \frac{1}{v} \frac{d\vec{R}}{dt}$  tangent to curve

$$\left| \frac{d\vec{R}}{ds} \right| = \frac{\left| \frac{d\vec{R}}{dt} \right|}{ds/dt} = 1$$

$$\Rightarrow \boxed{\vec{T}(s) = \frac{d\vec{R}}{ds}} \text{ : unit tangent vector}$$

eg Helix,  $\vec{R}(t) = p \sin t \hat{i} + p \cos t \hat{j} + at \hat{k}$   $0 \leq t \leq 4\pi$   
parametric form (Note:  $\hat{e}_1 = \hat{j}, \hat{e}_2 = \hat{i}, \hat{e}_3 = -\hat{k}$  (right-handed))  
[Equivalent non-parametric form:  $x = p \sin \frac{z}{a}, y = p \cos \frac{z}{a}$ ]  $p, a > 0$ : left-handed helix)

$$\vec{v}(t) = \vec{R}'(t) = p \cos t \hat{i} - p \sin t \hat{j} + a \hat{k}$$
$$v = |\vec{R}'(t)| = (p^2 \cos^2 t + p^2 \sin^2 t + a^2)^{1/2} = (p^2 + a^2)^{1/2}$$

$$s = \int_0^t |\vec{R}'(t)| dt = \int_0^t (p^2 + a^2)^{1/2} dt = (p^2 + a^2)^{1/2} t$$

eg  $\vec{R}(t) = 3 \sin t \hat{i} + 3 \cos t \hat{j} + 4t \hat{k}$   $p=3, a=4$   
 $\Rightarrow s = (3^2 + 4^2)^{1/2} t = 5t, \quad t = s/5$

$$\Rightarrow \vec{R}(s) = 3 \sin \frac{s}{5} \hat{i} + 3 \cos \frac{s}{5} \hat{j} + \frac{4}{5} s \hat{k}, \quad 0 \leq s \leq 20\pi$$

alternative parametrization

$$\vec{T}(s) = \vec{R}'(s) = \frac{3}{5} \cos \frac{s}{5} \hat{i} - \frac{3}{5} \sin \frac{s}{5} \hat{j} + \frac{4}{5} \hat{k}, \quad |\vec{T}| = 1.$$

## Acceleration and Curvature

Acceleration = rate of change of velocity :

Motivation for definitions, results: due to change in magnitude and/or direction

We know:

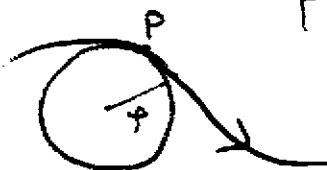
Change in magnitude only: motion along straight line

Tangential  $|\vec{a}| = \left| \frac{d\vec{v}}{dt} \right| = \frac{d}{dt} |\vec{v}| = \frac{d^2s}{dt^2}$

Change in direction only: motion along circle, radius  $\rho$

Centripetal  $|\vec{a}| = \frac{|\vec{v}|^2}{\rho} = \frac{1}{\rho} \left( \frac{ds}{dt} \right)^2 = \frac{v^2}{\rho}$

For general motion, acceleration is due to changes in both magnitude and direction.



For general curves: - what is  $\rho$ ?

Osculating circle best approximates curve at a given point P, tangent at P.

Radius  $\rho$  of osculating circle = radius of curvature  
Smaller radius  $\Leftrightarrow$  larger curvature

Characterizing curvature:

Unit tangent  $\vec{T} = \frac{\vec{R}'}{|\vec{R}'|} = \frac{1}{v} \frac{d\vec{R}}{ds}$ : fixed magnitude  $|\vec{T}| = 1$

$\Rightarrow \frac{d\vec{T}}{ds}$  due only to changes in direction

Greater curvature  $\Rightarrow$  greater change in  $\vec{T}$

$\Rightarrow$  define

Curvature  $k = \left| \frac{d\vec{T}}{ds} \right| = \frac{|d\vec{T}/dt|}{ds/dt} = \left| \frac{d^2\vec{R}}{ds^2} \right|$

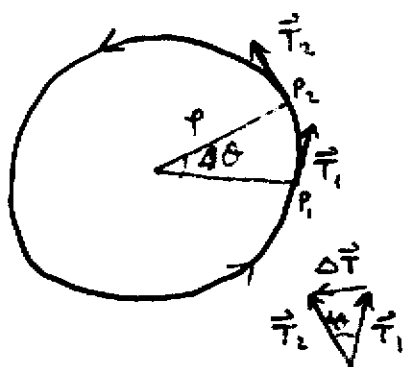
intrinsic properties of curves, independent of parametrization

Radius of curvature  $\rho = \frac{1}{k} = \frac{1}{|d\vec{T}/ds|}$   
(for  $k \neq 0$ )

$\uparrow$   
rate of change of  $\vec{T}$   
w.r.t. arc length



Check that the definitions agree with the usual radius for motion on a circle:



Arc Length from  $P_1$  to  $P_2$ :  $\Delta s \approx r \Delta \theta$

Magnitude of change in unit tangent

$$|\Delta \vec{T}| = |\vec{T}_2 - \vec{T}_1| \approx \frac{|\vec{T}_1|}{1} \Delta \theta = \Delta \theta$$

$$\Rightarrow \left| \frac{\Delta \vec{T}}{\Delta s} \right| \approx \frac{\Delta \theta}{\Delta s} \approx \frac{\Delta \theta}{r \Delta \theta} = \frac{1}{r}$$

limit as  $\Delta s \rightarrow 0$ :  $k = \left| \frac{d\vec{T}}{ds} \right| = \frac{1}{r}$  - agrees with above definition.

$\vec{T}$  has constant magnitude  
 $\Rightarrow$  either  $\frac{d\vec{T}}{dt} = \vec{0}$  or  $\frac{d\vec{T}}{dt} \perp \vec{T}$   $\left[ \begin{array}{l} \vec{T} \cdot \vec{T} = 1 \\ \frac{d}{dt} \vec{T} \cdot \vec{T} = 2\vec{T} \cdot \frac{d\vec{T}}{dt} = 0 \end{array} \right.$

If  $\frac{d\vec{T}}{dt} \neq 0$ , define the unit vector

Principal normal  $\vec{N} = \frac{\frac{d\vec{T}}{dt}}{\left| \frac{d\vec{T}}{dt} \right|} = \frac{\frac{d\vec{T}/ds}{\left| \frac{d\vec{T}}{ds} \right|}} = \frac{1}{k} \frac{d\vec{T}}{ds}$   
chain rule

Note:  $\vec{T} \cdot \frac{d\vec{T}}{dt} = 0 \Rightarrow \vec{T} \cdot \vec{N} = 0$  ie  $\vec{N} \perp \vec{T}$



$\vec{N}$ : a unit vector pointing to centre of osculating circle

Osculating plane: the plane containing  $\vec{T}$  and  $\vec{N}$

From the definition:  $\frac{d\vec{T}}{ds} = k \vec{N}$

ie  $\vec{T}$  turns in the direction  $\vec{N}$  at rate  $k$  per unit length of arc

Smooth arc  $\vec{R}(t) = x(t)\hat{i} + y(t)\hat{j} + z(t)\hat{k}$  (Assume  $\vec{R}(t) \in C^2$   
ie  $R_i(t)$  is twice continuously differentiable)

velocity  $\Rightarrow \vec{v}(t) = \vec{R}'(t) = \frac{dx}{dt}\hat{i} + \frac{dy}{dt}\hat{j} + \frac{dz}{dt}\hat{k}$

speed  $v(t) = |\vec{v}(t)| = \sqrt{\left(\frac{dx}{dt}\right)^2 + \left(\frac{dy}{dt}\right)^2 + \left(\frac{dz}{dt}\right)^2} = \frac{ds}{dt}$

acceleration  $\vec{a}(t) = \frac{d\vec{v}}{dt} = \vec{R}''(t) = \frac{d^2x}{dt^2}\hat{i} + \frac{d^2y}{dt^2}\hat{j} + \frac{d^2z}{dt^2}\hat{k}$

Derive an expression for  $\vec{a}$  in terms of  $\vec{T}$ ,  $\vec{N}$ :

Note  $\vec{v}(t) = \frac{d\vec{R}}{dt} = \frac{d\vec{R}}{ds} \frac{ds}{dt} \left[ = v \frac{d\vec{R}}{ds} \right] = \frac{ds}{dt} \vec{T}$

$$\begin{aligned} \Rightarrow \vec{a}(t) &= \frac{d\vec{v}}{dt} = \frac{d}{dt} \left( \frac{ds}{dt} \vec{T} \right) \quad \text{product rule} \\ &= \frac{d^2s}{dt^2} \vec{T} + \frac{ds}{dt} \frac{d\vec{T}}{dt} \\ &= \frac{d^2s}{dt^2} \vec{T} + \frac{ds}{dt} \frac{d\vec{T}}{ds} \frac{ds}{dt} \quad \frac{d\vec{T}}{ds} = \kappa \vec{N} \\ &= \frac{d^2s}{dt^2} \vec{T} + \kappa \left( \frac{ds}{dt} \right)^2 \vec{N} \end{aligned}$$

$\vec{a}$  lies in the osculating plane  $\rightarrow$

$$\vec{a} = a_{\parallel} \vec{T} + a_{\perp} \vec{N}$$

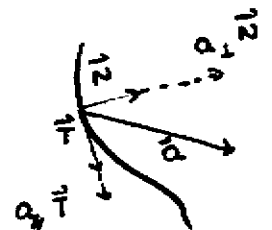
where (as expected)

$$a_{\parallel} = a_t = \frac{d^2s}{dt^2}$$

tangential acceleration

$$a_{\perp} = a_n = \kappa \left( \frac{ds}{dt} \right)^2 = \kappa v^2 = \frac{v^2}{\rho}$$

centripetal acceleration



Since  $\vec{T} \perp \vec{N}$ , we have (Pythagoras)

$$a^2 = |\vec{a}|^2 = a_{\parallel}^2 + a_{\perp}^2$$

• to find  $a$ , compute  $\frac{d^2\vec{R}}{dt^2}$ , calculate the magnitude

• to find  $a_{\parallel}$ , find  $v = \left| \frac{d\vec{R}}{dt} \right| = \frac{ds}{dt}$ , then differentiate wrt  $t$

• then find  $a_{\perp}$  from  $a_{\perp}^2 = a^2 - a_{\parallel}^2$

$$\text{eg } \vec{R}(t) = \underbrace{(s \sin t - t \cos t)}_{x(t)} \hat{i} + \underbrace{(t \cos t + s \sin t)}_{y(t)} \hat{j} + \underbrace{t^2}_{z(t)} \hat{k}$$

$$\frac{d\vec{R}}{dt} = (t \sin t) \hat{i} + (t \cos t) \hat{j} + 2t \hat{k} = \vec{v}(t)$$

$$\frac{d^2\vec{R}}{dt^2} = (t \cos t + \sin t) \hat{i} + (-t \sin t + \cos t) \hat{j} + 2 \hat{k} = \vec{a}(t)$$

$$\text{Speed: } v(t) = \left| \frac{d\vec{R}}{dt} \right| = \frac{ds}{dt} = [t^2 \sin^2 t + t^2 \cos^2 t + 4t^2]^{1/2} = (5t^2)^{1/2} = \sqrt{5}t$$

$$\text{Tangential acceleration: } a_t = \frac{d^2s}{dt^2} = \sqrt{5}$$

$$\begin{aligned} \text{Acceleration: } a^2 &= (t \cos t + \sin t)^2 + (-t \sin t + \cos t)^2 + 2^2 \\ &= t^2 \cos^2 t + 2t \sin t \cos t + \sin^2 t + t^2 \sin^2 t - 2t \cos t \sin t \\ &\quad + \cos^2 t + 4 \\ &= t^2 + 1 + 4 = t^2 + 5 \Rightarrow a = \sqrt{t^2 + 5} \end{aligned}$$

$$\text{Centripetal acceleration: } a_c^2 = a^2 - a_t^2 = t^2 + 5 - 5 = t^2 \Rightarrow a_c = t$$

$$\text{Curvature: } a_c = kv^2 \Rightarrow k = a_c/v^2 = t/5t^2 = 1/5t$$

$$\text{Radius of curvature: } \rho = 1/k = 5t$$

$$\text{Tangent vector: } \vec{T}(t) = \frac{\vec{R}'(t)}{v(t)} = \frac{1}{\sqrt{5}} \sin t \hat{i} + \frac{1}{\sqrt{5}} \cos t \hat{j} + \frac{2}{\sqrt{5}} \hat{k}$$

$$\text{Principal normal: } \vec{N}(t) = \frac{d\vec{T}/dt}{|d\vec{T}/dt|} = \frac{\frac{1}{\sqrt{5}} \cos t \hat{i} - \frac{1}{\sqrt{5}} \sin t \hat{j}}{1/\sqrt{5}} = \cos t \hat{i} - \sin t \hat{j}$$

$$\text{Arc length: } s(t) = \int_0^t \sqrt{5} t' dt' = \frac{\sqrt{5}}{2} t^2 \Rightarrow t = \frac{\sqrt{2s}}{5^{1/4}}$$

Another formula for curvature:

$$\vec{R}' = v \vec{T}, \quad \vec{R}'' = \frac{d^2s}{dt^2} \vec{T} + kv^2 \vec{N}$$

$$\Rightarrow \vec{R}' \times \vec{R}'' = v \vec{T} \times \left[ \frac{d^2s}{dt^2} \vec{T} + kv^2 \vec{N} \right] = kv^3 \vec{T} \times \vec{N}$$

$$\left[ \begin{array}{l} \text{since } \vec{T} = \frac{\vec{R}'}{|\vec{R}'|} = \frac{\vec{R}'}{v} \\ \vec{a} = a_t \vec{T} + a_c \vec{N} \\ \uparrow \quad \quad \uparrow \\ \frac{d^2s}{dt^2} \quad kv^2 \end{array} \right.$$

Define

$$\text{Binormal } \boxed{\vec{B} = \vec{T} \times \vec{N}}$$

since  $|\vec{T}| = |\vec{N}| = 1$ ,  $\vec{T} \perp \vec{N}$ ,  
we have  $|\vec{B}| = |\vec{T} \times \vec{N}| = 1$

$$\Rightarrow \vec{R}' \times \vec{R}'' = kv^3 \vec{T} \times \vec{N} = kv^3 \vec{B}$$

$$\Rightarrow |\vec{R}' \times \vec{R}''| = kv^3 |\vec{B}| = k |\vec{R}'|^3$$

$$\Rightarrow \boxed{k = \frac{|\vec{R}' \times \vec{R}''|}{|\vec{R}'|^3}} = \frac{|\vec{v} \times \vec{a}|}{v^3} \quad \text{and } a_c = kv^2 = \frac{|\vec{R}' \times \vec{R}''|}{|\vec{R}'|}$$

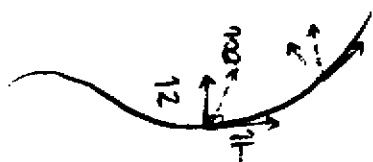
## Frenet Formulas (Frenet - Serret) - some differential geometry...

Unit tangent  $\vec{T} = \frac{d\vec{R}}{ds}$  , curvature  $k = \left| \frac{d\vec{T}}{ds} \right| = \left| \frac{d^2\vec{R}}{ds^2} \right|$

Principal normal  $\vec{N} = \frac{1}{k} \frac{d\vec{T}}{ds} \Rightarrow |\vec{T}| = 1, |\vec{N}| = 1, \vec{T} \cdot \vec{N} = 0$

Binormal  $\vec{B} = \vec{T} \times \vec{N} \Rightarrow |\vec{B}| = 1, \vec{B} \cdot \vec{T} = \vec{B} \cdot \vec{N} = 0$

$\vec{T}, \vec{N}, \vec{B}$  form a right-handed system of orthonormal vectors: Frenet frame



As we move along the curve, the triad of unit vectors moves and rotates

Motivation for Frenet formulas: (not a proof)

$$\frac{d\vec{T}}{ds} = \underbrace{k}_{\text{①}} \vec{N} \quad \left. \begin{array}{l} \vec{T} \text{ turns towards } \vec{N} \text{ at rate } k \text{ ①} \\ \Rightarrow \vec{N} \text{ turns towards } -\vec{T} \text{ at rate } k \text{ ②} \end{array} \right\} \begin{array}{l} \text{since} \\ \vec{T} \perp \vec{N} \\ \text{always} \end{array}$$

Also,  $\vec{N}$  could rotate about  $\vec{T}$ : tilt of osculating plane instantaneously containing curve (twist)

→  $\tau$ : torsion: measures rate at which curve twists  
ie  $\frac{d\vec{N}}{ds}$  could have a component  $\perp \vec{T}$   
(and  $\perp \vec{N}$ :  $\vec{N} \cdot \vec{N} = 1 \Rightarrow \vec{N} \cdot \frac{d\vec{N}}{ds} = 0$ ) ie  $\parallel \vec{B}$  ③

Thus

$$\frac{d\vec{N}}{ds} = \underbrace{-k}_{\text{②}} \vec{T} + \underbrace{\tau}_{\text{③}} \vec{B} \quad \text{- twist}$$

$\vec{N}$  turns towards  $\vec{B}$  at rate  $\tau$  ③  $\Rightarrow \vec{B}$  turns towards  $\vec{N}$  at rate  $-\tau$  ④

$\vec{B}$  does not turn towards  $\vec{T}$ , since  $\vec{T}$  turns only towards  $\vec{N}$ , by definition

Thus

$$\frac{d\vec{B}}{ds} = \underbrace{-\tau}_{\text{④}} \vec{N}$$

Derivation of Frenet formulas:

$$1. \quad \frac{d\vec{T}}{ds} = k \vec{N}, \text{ and } k = \left| \frac{d\vec{T}}{ds} \right| \geq 0$$

(definition). curvature

$$\begin{aligned} \frac{d\vec{T}}{ds} &= k \vec{N} \\ \frac{d\vec{N}}{ds} &= -k \vec{T} + \tau \vec{B} \\ \frac{d\vec{B}}{ds} &= -\tau \vec{N} \end{aligned}$$

$$2. \quad \vec{B} = \vec{T} \times \vec{N} \Rightarrow \vec{B} \cdot \vec{B} = |\vec{B}|^2 = 1$$

$$\Rightarrow \vec{B} \cdot \frac{d\vec{B}}{ds} = \frac{1}{2} \frac{d}{ds} (\vec{B} \cdot \vec{B}) = 0 \Rightarrow \frac{d\vec{B}}{ds} \perp \vec{B}$$

$$\begin{aligned} \frac{d\vec{B}}{ds} &= \frac{d}{ds} (\vec{T} \times \vec{N}) = \frac{d\vec{T}}{ds} \times \vec{N} + \vec{T} \times \frac{d\vec{N}}{ds} = k \vec{N} \times \vec{N} + \vec{T} \times \frac{d\vec{N}}{ds} \\ &= \vec{T} \times \frac{d\vec{N}}{ds}, \text{ so } \frac{d\vec{B}}{ds} \perp \vec{T}. \end{aligned}$$

Since  $\frac{d\vec{B}}{ds} \perp \vec{B}$ , and  $\perp \vec{T}$ , we must have  $\frac{d\vec{B}}{ds} \parallel \vec{N}$

Thus we can write  $\frac{d\vec{B}}{ds} = -\tau \vec{N}$  ← this formula defines

$$\text{and } \tau = -\vec{N} \cdot \frac{d\vec{B}}{ds}$$

the torsion  $\tau$ .

scalar ("amount of twist")

3.  $\vec{T}, \vec{N}, \vec{B}$  form a right-handed system

$$\Rightarrow \vec{N} = \vec{B} \times \vec{T}$$

$$\Rightarrow \frac{d\vec{N}}{ds} = \frac{d\vec{B}}{ds} \times \vec{T} + \vec{B} \times \frac{d\vec{T}}{ds} = \underbrace{(-\tau) \vec{N}}_{-\vec{B}} \times \vec{T} + k \underbrace{\vec{B} \times \vec{N}}_{-\vec{T}}$$

$$\Rightarrow \frac{d\vec{N}}{ds} = -k \vec{T} + \tau \vec{B}$$

Notes:

- For planar motion,  $\vec{T}, \vec{N}$  lie in the plane of the curve,  $\vec{B}$  is a constant unit vector  $\perp$  plane, torsion  $\tau = 0$  (no twist)

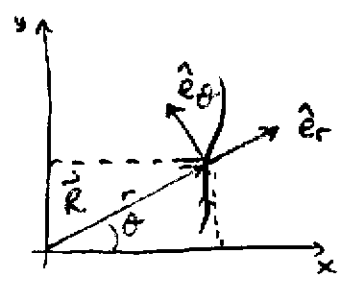
$k = k(s), \tau = \tau(s)$   
intrinsic equations  
of curve

- Two curves with identical  $k(s), \tau(s)$  are congruent (can coincide after a rigid motion)  $\Rightarrow k(s), \tau(s)$  describe curve

• Formulas for torsion:  $[\vec{r}', \vec{r}'', \vec{r}''']$

$$\tau = \frac{[\vec{r}', \vec{r}'', \vec{r}''']}{|\vec{r}' \times \vec{r}''|^2} = \frac{1}{k^2} \left[ \frac{d\vec{r}}{ds}, \frac{d^2\vec{r}}{ds^2}, \frac{d^3\vec{r}}{ds^3} \right]$$

# Planar Motion in Polar Coordinates



$$\left. \begin{aligned} x &= r \cos \theta \\ y &= r \sin \theta \end{aligned} \right\} \begin{aligned} r &\geq 0 \\ -\pi < \theta &\leq \pi \end{aligned}$$

$$\Rightarrow r = (x^2 + y^2)^{1/2}$$

$$\tan \theta = \frac{y}{x} \Rightarrow \theta = \arctan \frac{y}{x} = \arcsin \frac{y}{(x^2 + y^2)^{1/2}}$$

avoid quadrant ambiguities

Radial and angular unit vectors  $\hat{e}_r, \hat{e}_\theta$  (or  $\hat{u}_r, \hat{u}_\theta$ )

vary from point to point

$$\left. \begin{aligned} \hat{e}_r &= \cos \theta \hat{i} + \sin \theta \hat{j} \\ \hat{e}_\theta &= -\sin \theta \hat{i} + \cos \theta \hat{j} \end{aligned} \right\} \begin{aligned} \text{Note } \hat{e}_r \cdot \hat{e}_r &= \hat{e}_\theta \cdot \hat{e}_\theta = 1 \\ \hat{e}_r \cdot \hat{e}_\theta &= 0 \end{aligned}$$

not constant!

$$\frac{\partial \hat{e}_r}{\partial r} = 0, \quad \frac{\partial \hat{e}_\theta}{\partial r} = 0, \quad \frac{\partial \hat{e}_r}{\partial \theta} = -\sin \theta \hat{i} + \cos \theta \hat{j} = \hat{e}_\theta, \quad \frac{\partial \hat{e}_\theta}{\partial \theta} = -\hat{e}_r$$

Chain rule  $\Rightarrow \frac{d\hat{e}_r}{dt} = \frac{\partial \hat{e}_r}{\partial r} \frac{dr}{dt} + \frac{\partial \hat{e}_r}{\partial \theta} \frac{d\theta}{dt} = \frac{d\theta}{dt} \hat{e}_\theta, \quad \frac{d\hat{e}_\theta}{dt} = -\frac{d\theta}{dt} \hat{e}_r$

Motion in polar coordinates:

$\vec{R}(t) = r \hat{e}_r$  position vector of a particle at  $(r, \theta)$

velocity  $\vec{v}(t) = \frac{d\vec{R}}{dt} = \frac{dr}{dt} \hat{e}_r + r \frac{d\hat{e}_r}{dt} = \frac{dr}{dt} \hat{e}_r + r \frac{d\theta}{dt} \hat{e}_\theta$

$$\Rightarrow \vec{v}(t) = \underbrace{\frac{dr}{dt}}_{\text{radial}} \hat{e}_r + r \underbrace{\frac{d\theta}{dt}}_{\text{transverse}} \hat{e}_\theta$$

acceleration  $\vec{a}(t) = \frac{d^2\vec{R}}{dt^2} = \frac{d^2r}{dt^2} \hat{e}_r + \frac{dr}{dt} \frac{d\hat{e}_r}{dt} + \frac{dr}{dt} \frac{d\theta}{dt} \hat{e}_\theta + r \frac{d^2\theta}{dt^2} \hat{e}_\theta + r \frac{d\theta}{dt} \frac{d\hat{e}_\theta}{dt}$

$$= \left[ \frac{d^2r}{dt^2} - r \frac{d\theta}{dt} \right] \hat{e}_r + \left[ r \frac{d^2\theta}{dt^2} + 2 \frac{dr}{dt} \frac{d\theta}{dt} \right] \hat{e}_\theta$$

Velocity  $\vec{v} = \frac{d\vec{r}}{dt} = \underbrace{\frac{dr}{dt} \hat{e}_r}_{\text{radial}} + r \underbrace{\frac{d\theta}{dt} \hat{e}_\theta}_{\text{tangential}}$

Acceleration  $\vec{a} = \frac{d^2\vec{r}}{dt^2} = \left[ \frac{d^2r}{dt^2} - r \left( \frac{d\theta}{dt} \right)^2 \right] \hat{e}_r + \left[ r \frac{d^2\theta}{dt^2} + 2 \frac{dr}{dt} \frac{d\theta}{dt} \right] \hat{e}_\theta$

$\frac{d^2r}{dt^2} \hat{e}_r$  : acceleration for pure radial motion

$r \frac{d^2\theta}{dt^2} \hat{e}_\theta$  : pure angular acceleration

$-r \left( \frac{d\theta}{dt} \right)^2 \hat{e}_r$  : centripetal acceleration (if  $r = \text{constant}$ )

$2 \frac{dr}{dt} \frac{d\theta}{dt} \hat{e}_\theta$  : related to Coriolis acceleration { due to change in direction of  $\hat{e}_r$ , & transverse component changes with  $r$  even if  $\frac{d\theta}{dt} = \text{const.}$

Special cases:

1.  $\theta = \text{constant}$ ,  $\vec{r} = r \hat{e}_r \leftarrow \text{constant}$

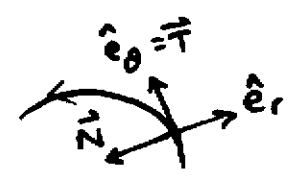
$\Rightarrow \vec{a} = \frac{d^2r}{dt^2} \hat{e}_r = \frac{d^2s}{dt^2} \hat{T}$  (since  $s=r$ )

- one-dimensional motion (straight line)

2.  $r = \text{constant}$

- circular motion

$\vec{a} = -r \left( \frac{d\theta}{dt} \right)^2 \hat{e}_r + r \frac{d^2\theta}{dt^2} \hat{e}_\theta$   
 $= r \left( \frac{d\theta}{dt} \right)^2 \vec{N} + r \frac{d^2\theta}{dt^2} \hat{T}$



[and  $s = r\theta \Rightarrow \frac{ds}{dt} = r \frac{d\theta}{dt}$ ,  $\frac{d^2s}{dt^2} = r \frac{d^2\theta}{dt^2}$ ]

$\Rightarrow \vec{a} = k \left( \frac{ds}{dt} \right)^2 \vec{N} + \frac{d^2s}{dt^2} \hat{T}$

[curvature  $k = \frac{1}{r}$ ]

Newton's 2<sup>nd</sup> Law:  $\vec{F} = m\vec{a}$

$$\vec{F} = F_r \hat{e}_r + F_\theta \hat{e}_\theta$$

write force i.b. radial, angular components

$$\Rightarrow F_r = m \frac{d^2 r}{dt^2} - m r \left( \frac{d\theta}{dt} \right)^2 = m \ddot{r} - m r \dot{\theta}^2$$

time derivative  
convention:  $\dot{r} \equiv \frac{dr}{dt}$

and

$$F_\theta = m r \frac{d^2 \theta}{dt^2} + 2 m r \frac{dr}{dt} \frac{d\theta}{dt}$$

$$\Rightarrow \underbrace{r F_\theta}_{\text{torque}} = m r^2 \frac{d^2 \theta}{dt^2} + 2 m r^2 \frac{dr}{dt} \frac{d\theta}{dt} = \frac{d}{dt} \left[ \underbrace{m r^2 \frac{d\theta}{dt}}_{\text{angular momentum}} \right]$$

Applied torque = time rate of change of angular momentum

In general: Newton's 2<sup>nd</sup> Law  $\vec{F} = m \frac{d\vec{v}}{dt}$

$$\text{Torque } \vec{\tau} = \vec{R} \times \vec{F} = \vec{R} \times m \frac{d\vec{v}}{dt} = \frac{d}{dt} (\vec{R} \times m\vec{v}) \quad \left( \text{since } \frac{d\vec{R}}{dt} \times m\vec{v} = \vec{0} \right)$$

$$= \frac{d\vec{L}}{dt}, \quad \text{where } \vec{L} = \vec{R} \times \vec{p} = \vec{R} \times (m\vec{v})$$

is the angular momentum

In this case  $\vec{\tau} = r \hat{e}_r \times (F_r \hat{e}_r + F_\theta \hat{e}_\theta) = r F_\theta \hat{e}_r \times \hat{e}_\theta = r F_\theta \hat{k}$

and  $\vec{L} = \vec{R} \times m\vec{v} = m r \hat{e}_r \times \left( \frac{dr}{dt} \hat{e}_r + r \frac{d\theta}{dt} \hat{e}_\theta \right) = m r^2 \frac{d\theta}{dt} \hat{k}$

If  $F_\theta = 0$  i.e.  $\vec{F} = F_r \hat{e}_r$  : radial (central) force

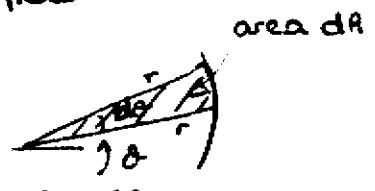
then  $\frac{d}{dt} \left[ m r^2 \frac{d\theta}{dt} \right] = 0 \Rightarrow m r^2 \frac{d\theta}{dt} = \text{const} = C$

angular momentum is constant in a central force field

Geometric interpretation:

$$dA = \frac{1}{2} r^2 d\theta$$

(infinitesimal area as radius vector rotates through angle  $d\theta$ )



$$\Rightarrow A(t) = \int_{t_1}^{t_2} \frac{1}{2} r^2 \frac{d\theta}{dt} dt \Rightarrow \frac{dA}{dt} = \frac{1}{2} r^2 \frac{d\theta}{dt} = \text{constant}$$

For a particle moving under a central force, equal areas are swept out in equal times: Kepler's 2<sup>nd</sup> Law of planetary motion